

Advanced Structural Geology, Fall 2022

Static stress changes-- Coulomb

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Static Stress Changes and the Triggering of Earthquakes

by Geoffrey C. P. King, Ross S. Stein, and Jian Lin

Key concepts:

- Source faults
- Receiver faults
- Optimally oriented faults
- Assume receiver faults are close to failure
- Triggering lag time is a problem

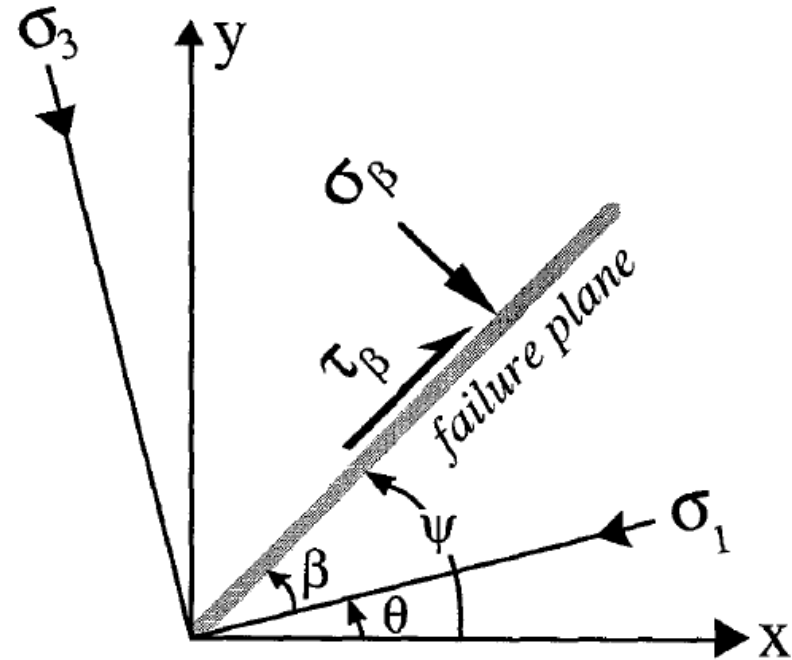
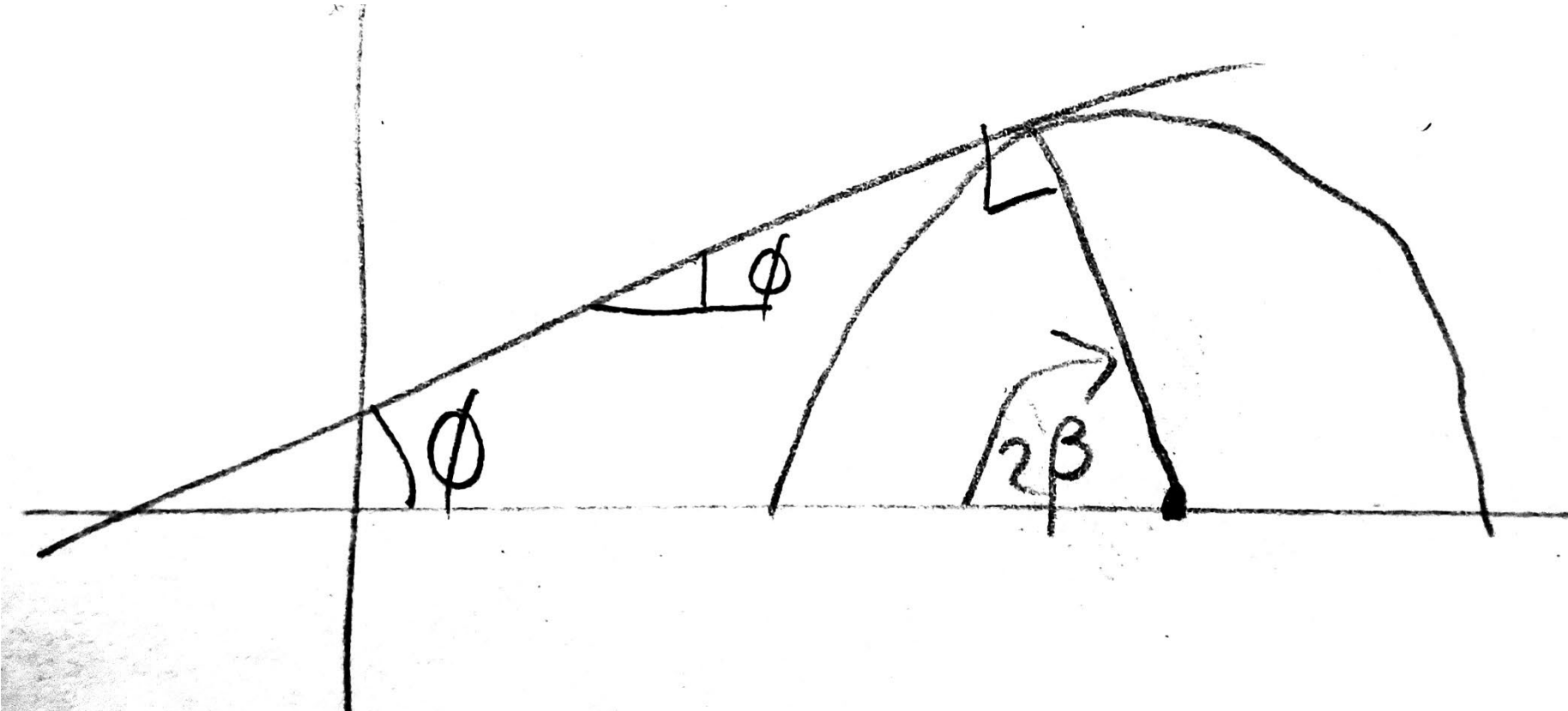


Figure 1. The axis system used for calculations of Coulomb stresses on optimum failure planes. Compression and right-lateral shear stress on the failure plane are taken as positive. The sign of τ_β is reversed for calculations of right-lateral Coulomb failure on specified failure planes.

Coulomb failure and optimal orientation of fault planes

Optimal orientation; $\phi + 90 + 2\beta = 180$



Change of coulomb stress on faults of specified orientation (see next lecture)

Given \mathbf{S} principal stress tensor with orientation $x'y'z'$

Rotate to N-S, E-W components

$$\mathbf{R} = \begin{bmatrix} l & l' & l'' \\ m & m' & m'' \\ n & n' & n'' \end{bmatrix} \text{ where } \begin{array}{l} l = x * x' \\ l' = x * y' \\ l'' = x * z' \\ \dots \end{array}$$

$$\mathbf{S}' = \mathbf{R}^T \mathbf{S} \mathbf{R}$$

Can change spatially
Remote: S_{ij}^r

Induced: S_{ij}^f

Total: $\mathbf{S} = \mathbf{S}_{ij}^r + \mathbf{S}_{ij}^f$

And given plane with normal vector direction cosines \mathbf{N}

Can change spatially

Traction $\mathbf{T} = \mathbf{S}' * \mathbf{N}$ (row and column multiplication)

$T = \sqrt{\mathbf{T}(1)^2 + \mathbf{T}(2)^2 + \mathbf{T}(3)^2}$ traction magnitude

$\sigma_n = \mathbf{T} \cdot \mathbf{N}$ dot product for normal traction magnitude

$\mathbf{B} = \mathbf{T} \times \mathbf{N}$ cross product for null vector

$B = \sqrt{\mathbf{B}(1)^2 + \mathbf{B}(2)^2 + \mathbf{B}(3)^2}$ \mathbf{B} magnitude

$\mathbf{B}_{normalized} = \mathbf{B}./B$ normalize for orientation if necessary

$\mathbf{T}_s = \mathbf{N} \times \mathbf{B}$ cross product for shear traction vector

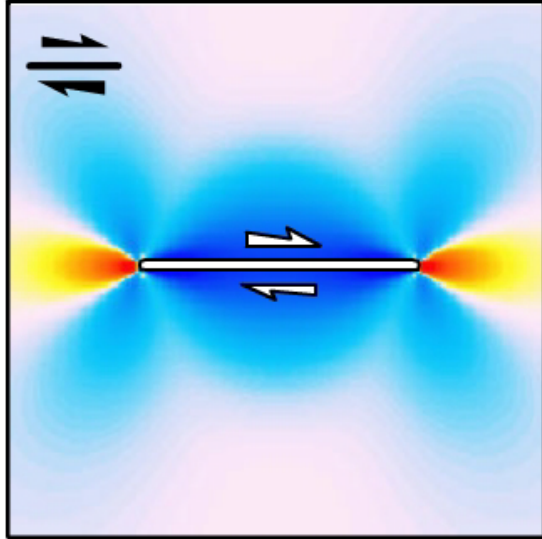
$\tau = \sqrt{\mathbf{T}_s(1)^2 + \mathbf{T}_s(2)^2 + \mathbf{T}_s(3)^2}$ shear traction magnitude

$\mathbf{T}_{snormalized} = \mathbf{T}_s./\tau$ normalize for shear traction orientation

Coulomb failure function: $\Delta\sigma_f = \Delta\tau - (\mu - P)\Delta\sigma_n$

How the Coulomb Stress Change is Calculated

Stress ■ Rise ■ Drop



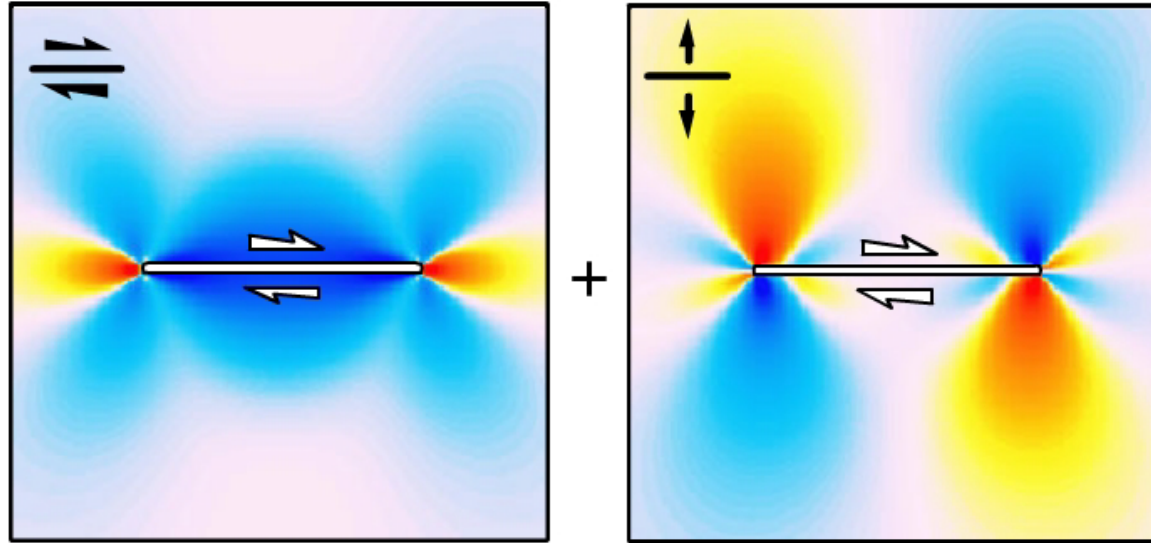
Shear stress
change

$$\Delta\tau_s$$

- Example calculation for faults parallel to master fault

How the Coulomb Stress Change is Calculated

Stress ■ Rise ■ Drop



Shear stress
change

$$\Delta\tau_s$$

+

Friction coefficient x
normal stress change

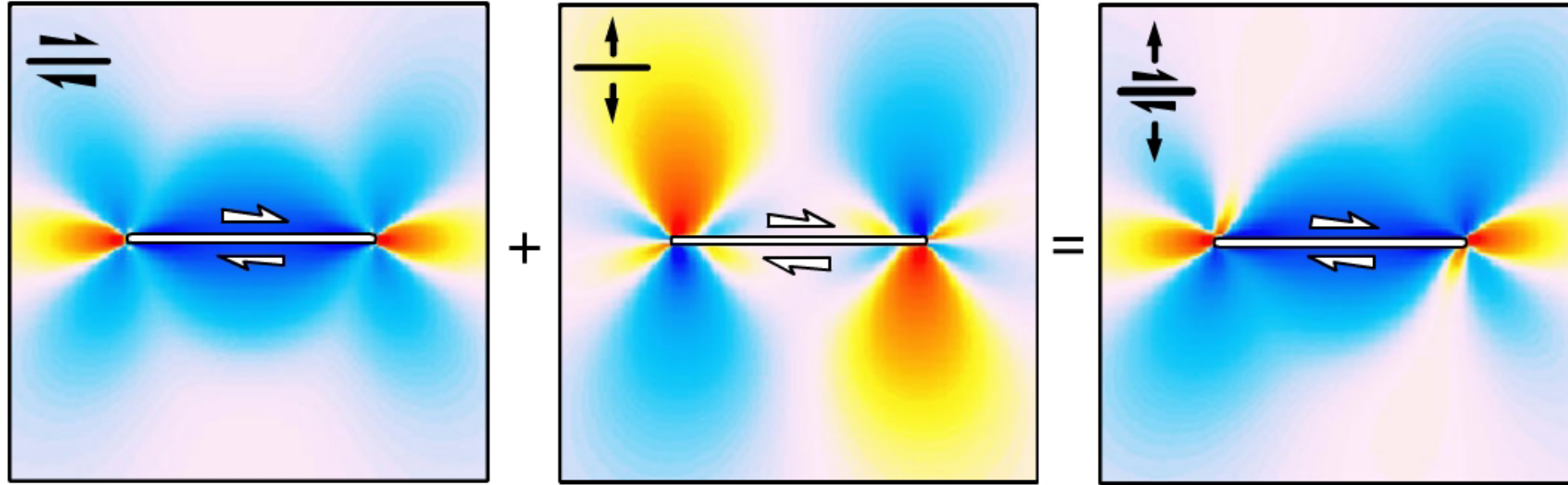
+

$$\mu' (\Delta\sigma_n)$$

- Example calculation for faults parallel to master fault

How the Coulomb Stress Change is Calculated

Stress ■ Rise ■ Drop



Shear stress
change

+

Friction coefficient x
normal stress change

=

Coulomb failure
stress change

$$\Delta\tau_s$$

+

$$\mu' (\Delta\sigma_n)$$

=

$$\Delta\sigma_f$$

- Example calculation for faults parallel to master fault

Change of coulomb stress on faults of optimal orientation

Mostly book keeping going from the stress tensor to the traction vector β is orientation from σ_1 of optimal failure plane

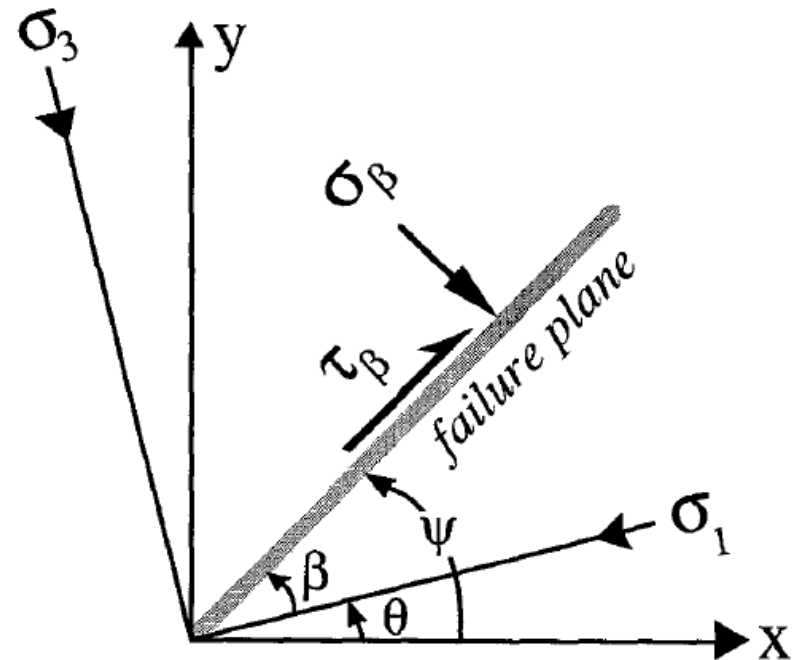


Figure 1. The axis system used for calculations of Coulomb stresses on optimum failure planes. Compression and right-lateral shear stress on the failure plane are taken as positive. The sign of τ_β is reversed for calculations of right-lateral Coulomb failure on specified failure planes.

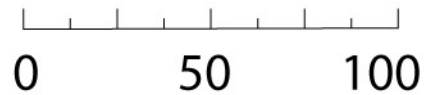
1986 $M=6.0$ North Palm Springs

Coulomb stress
imparted by
mainshocks

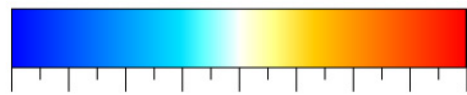
Source fault



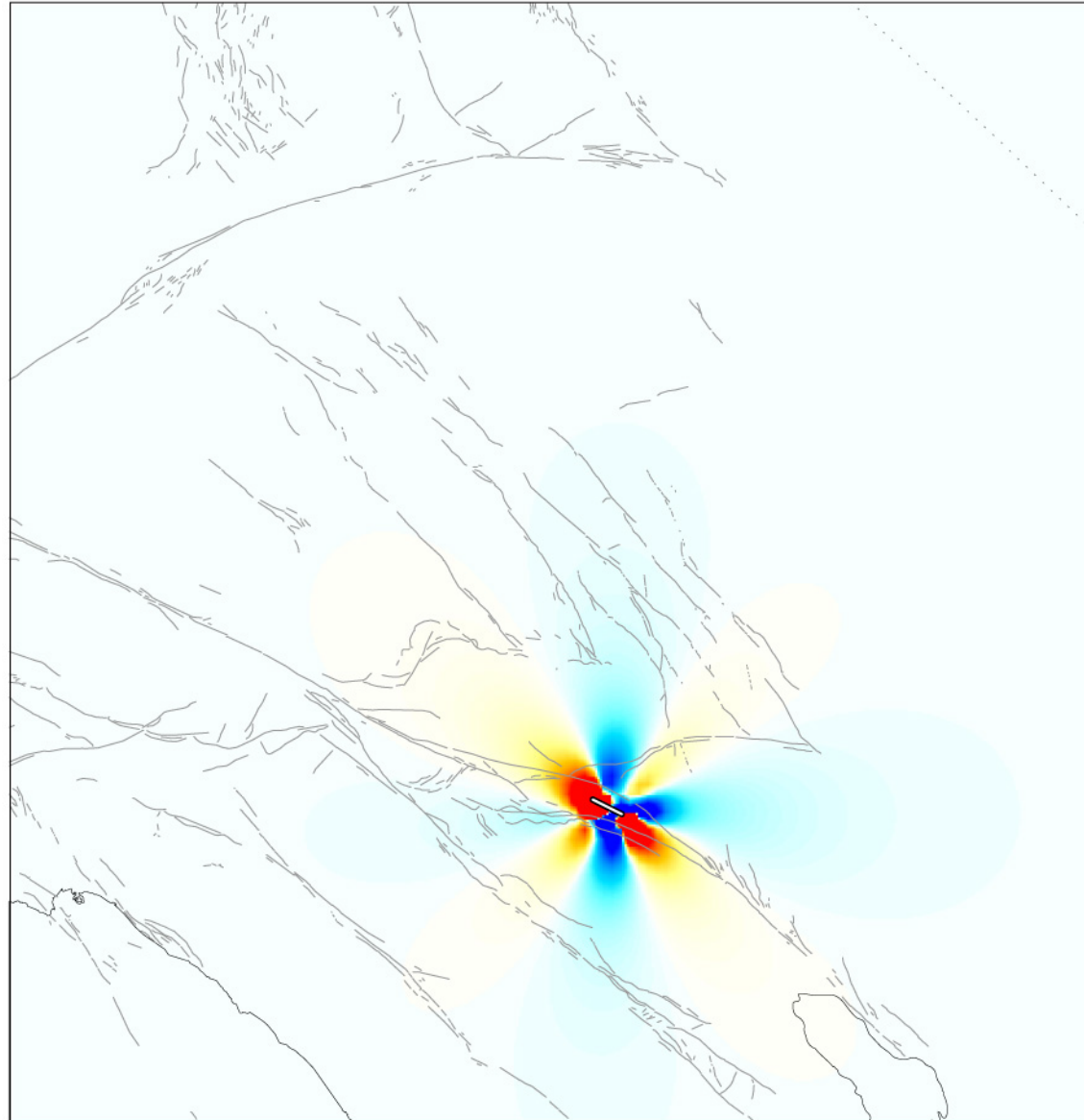
Distance (km)



Coulomb stress
change (bars)



-1.0 -0.5 0.0 0.5 1.0



from Todal et al (JGR, 2005)

1992 $M=6.2$ Joshua Tree

Coulomb stress
imparted by
mainshocks

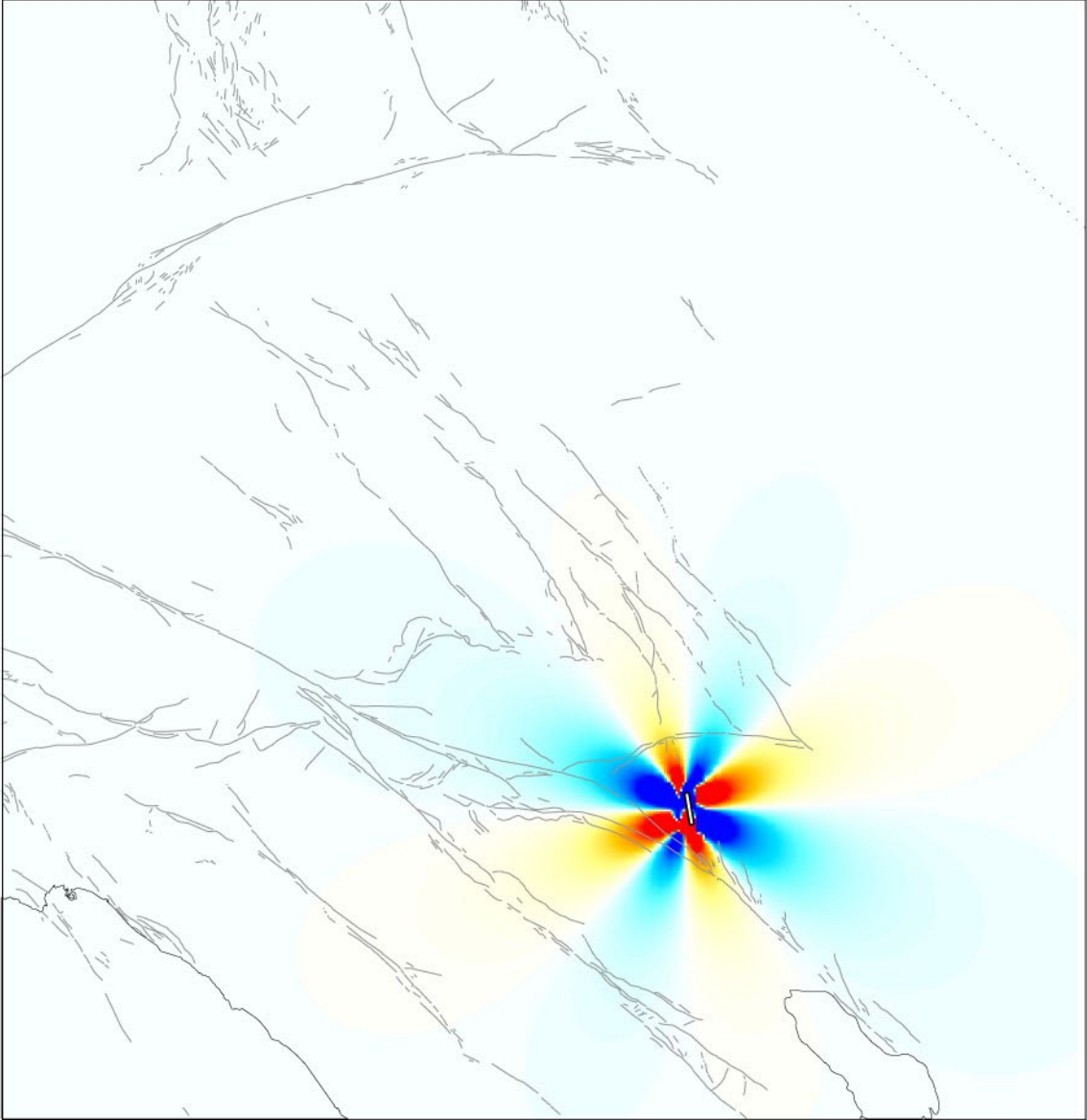
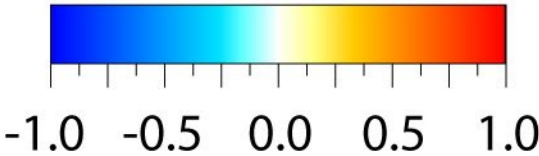
Source fault



Distance (km)



Coulomb stress
change (bars)



from Todal et al (JGR, 2005)

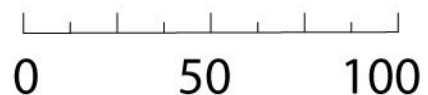
1992 $M=7.4$ Landers

Coulomb stress
imparted by
mainshocks

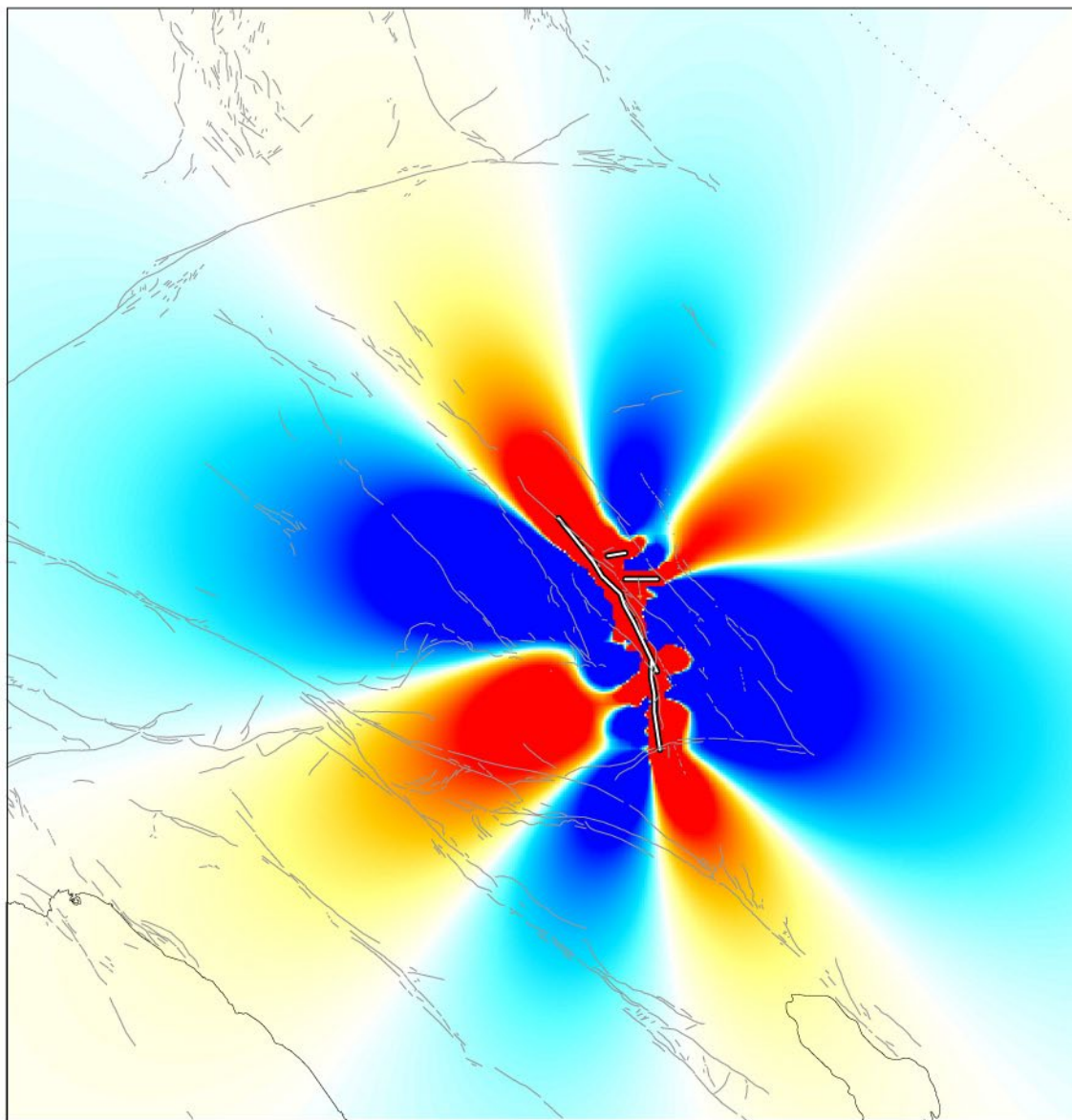
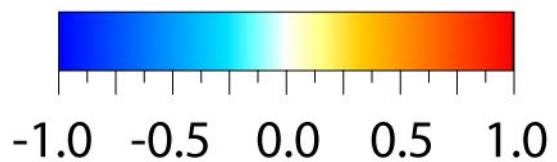
Source fault



Distance (km)



Coulomb stress
change (bars)



from Todal et al (JGR, 2005)

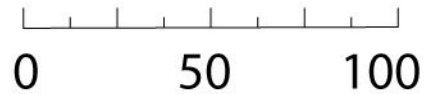
1992 $M=6.5$ Big Bear

Coulomb stress
imparted by
mainshocks

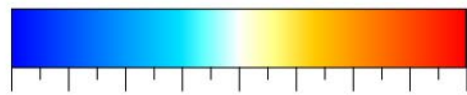
Source fault



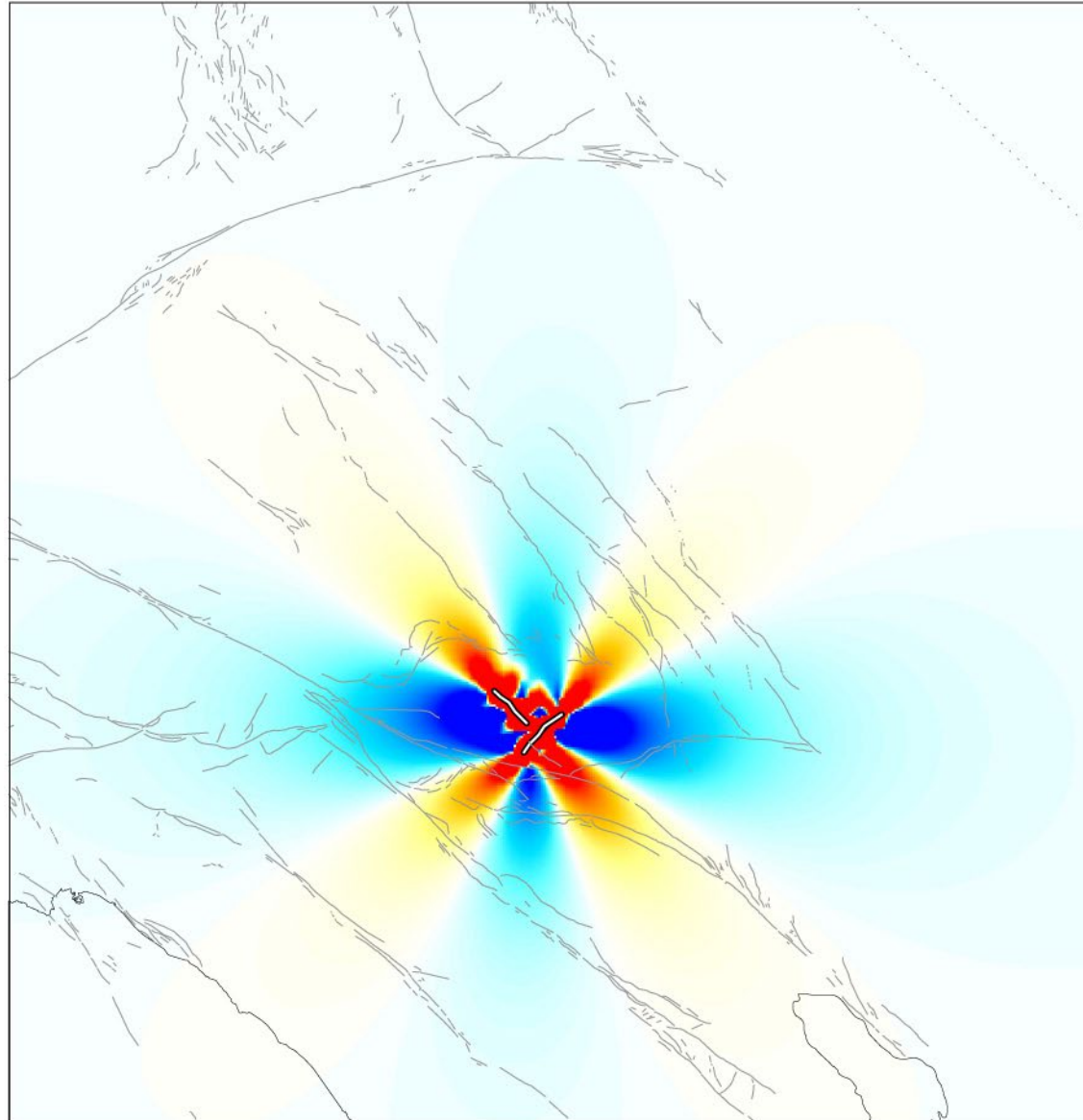
Distance (km)



Coulomb stress
change (bars)



-1.0 -0.5 0.0 0.5 1.0



from Todal et al (JGR, 2005)

1999 **M**=7.1 Hector Mine

Coulomb stress
imparted by
mainshocks

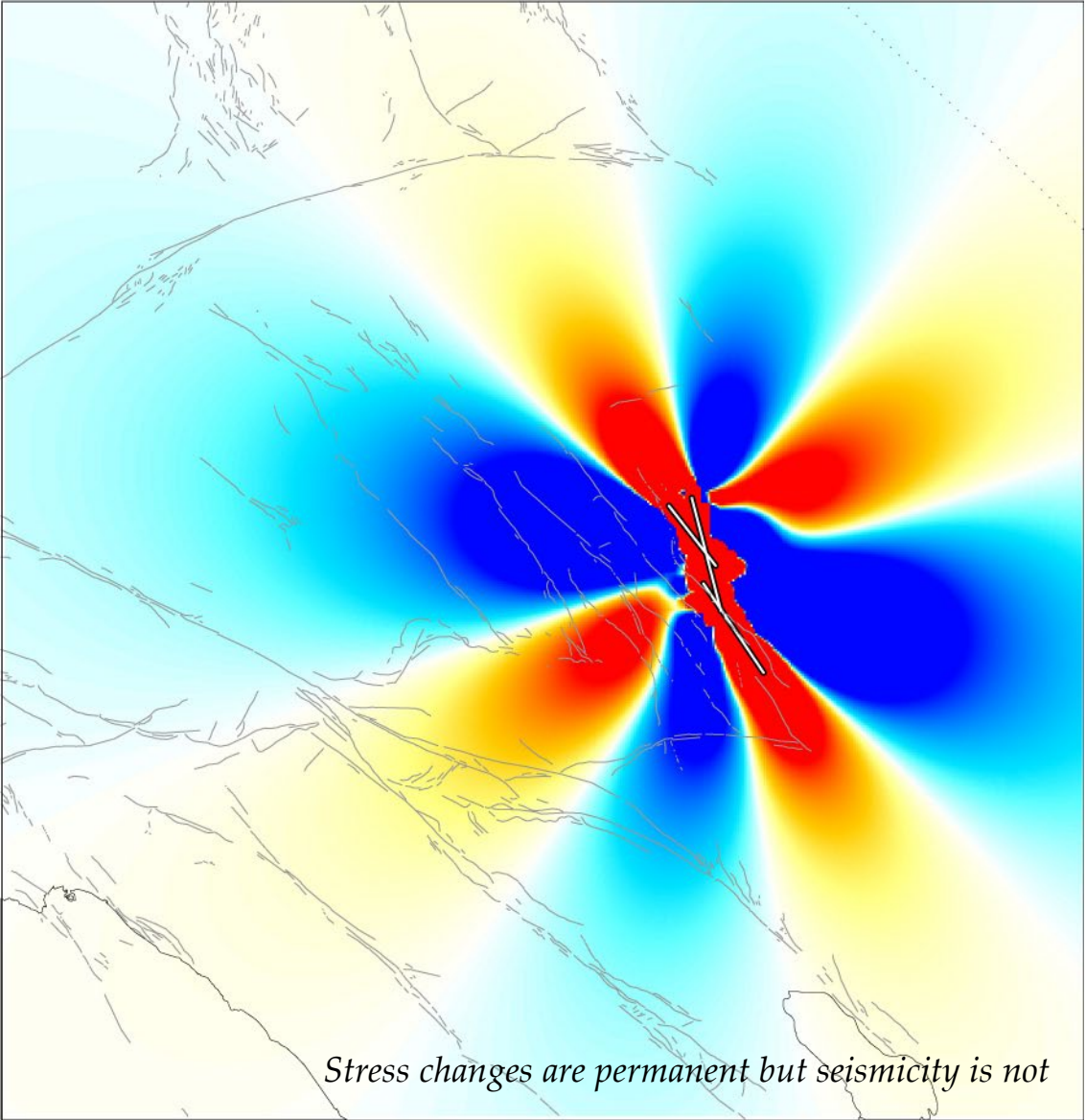
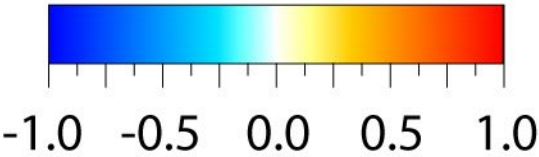
Source fault



Distance (km)

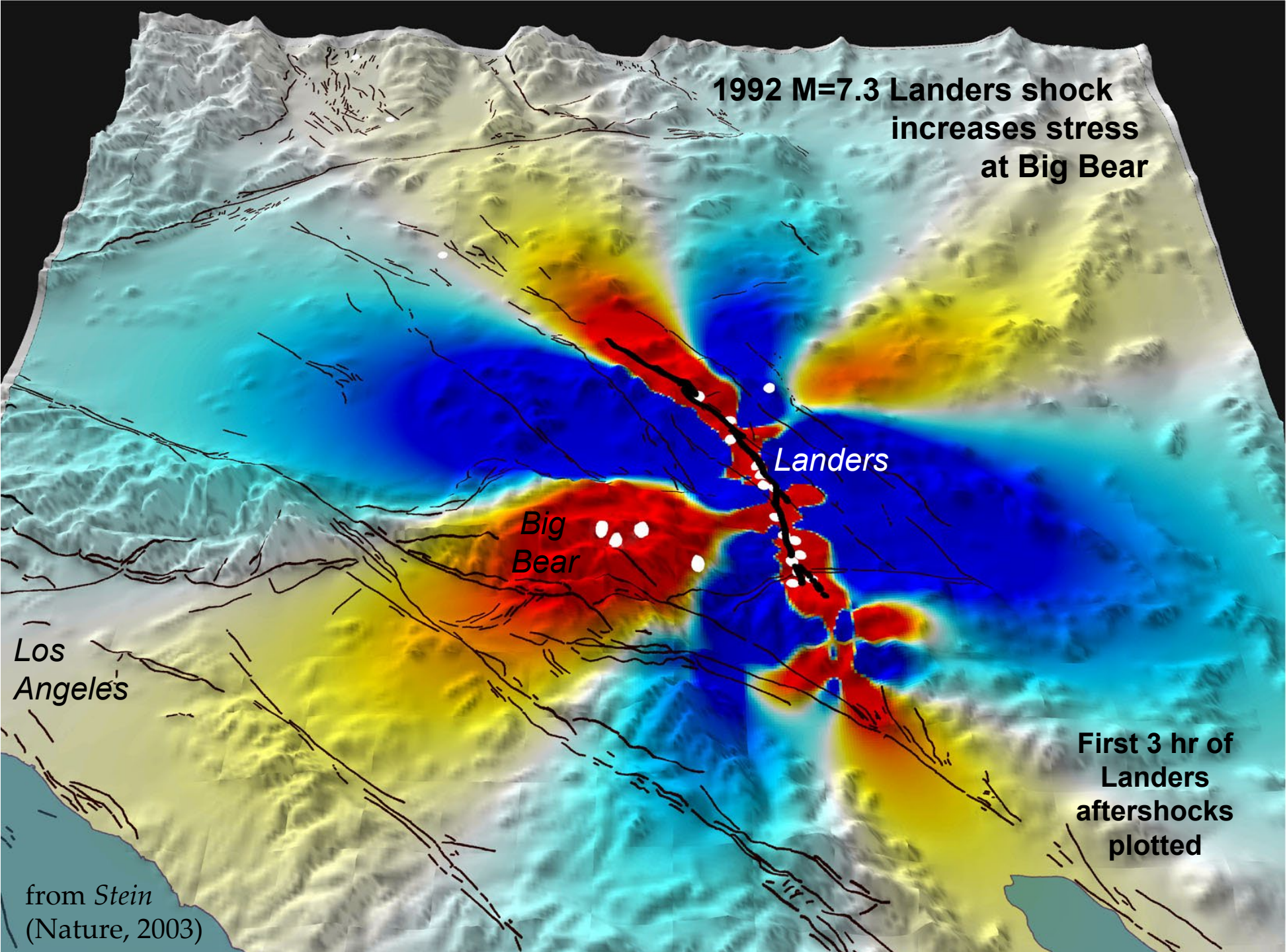


Coulomb stress
change (bars)



Stress changes are permanent but seismicity is not

from Todal et al (JGR, 2005)



1992 M=7.3 Landers shock
increases stress
at Big Bear

Landers

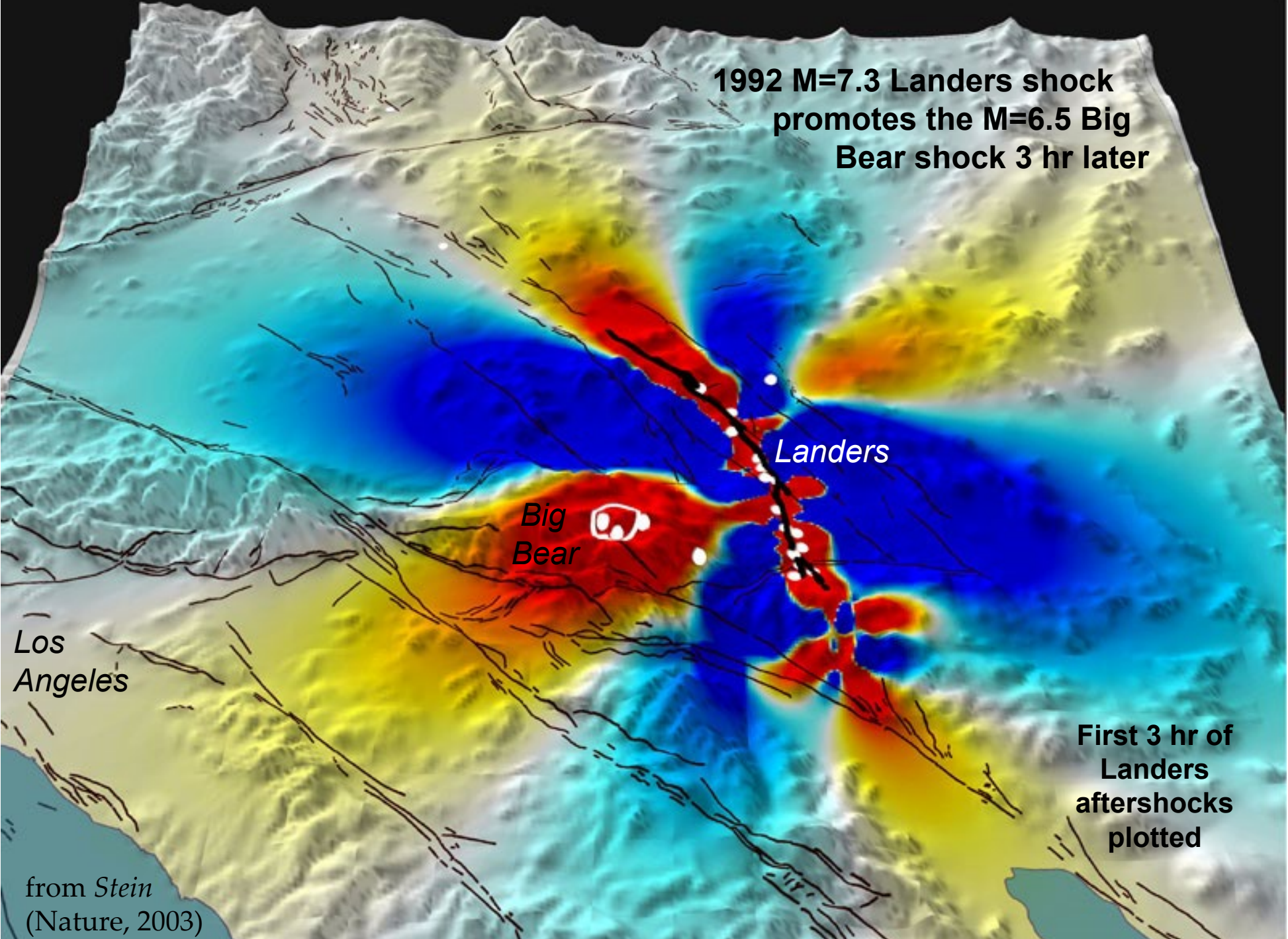
Big
Bear

Los
Angeles

First 3 hr of
Landers
aftershocks
plotted

from Stein
(Nature, 2003)

**1992 M=7.3 Landers shock
promotes the M=6.5 Big
Bear shock 3 hr later**



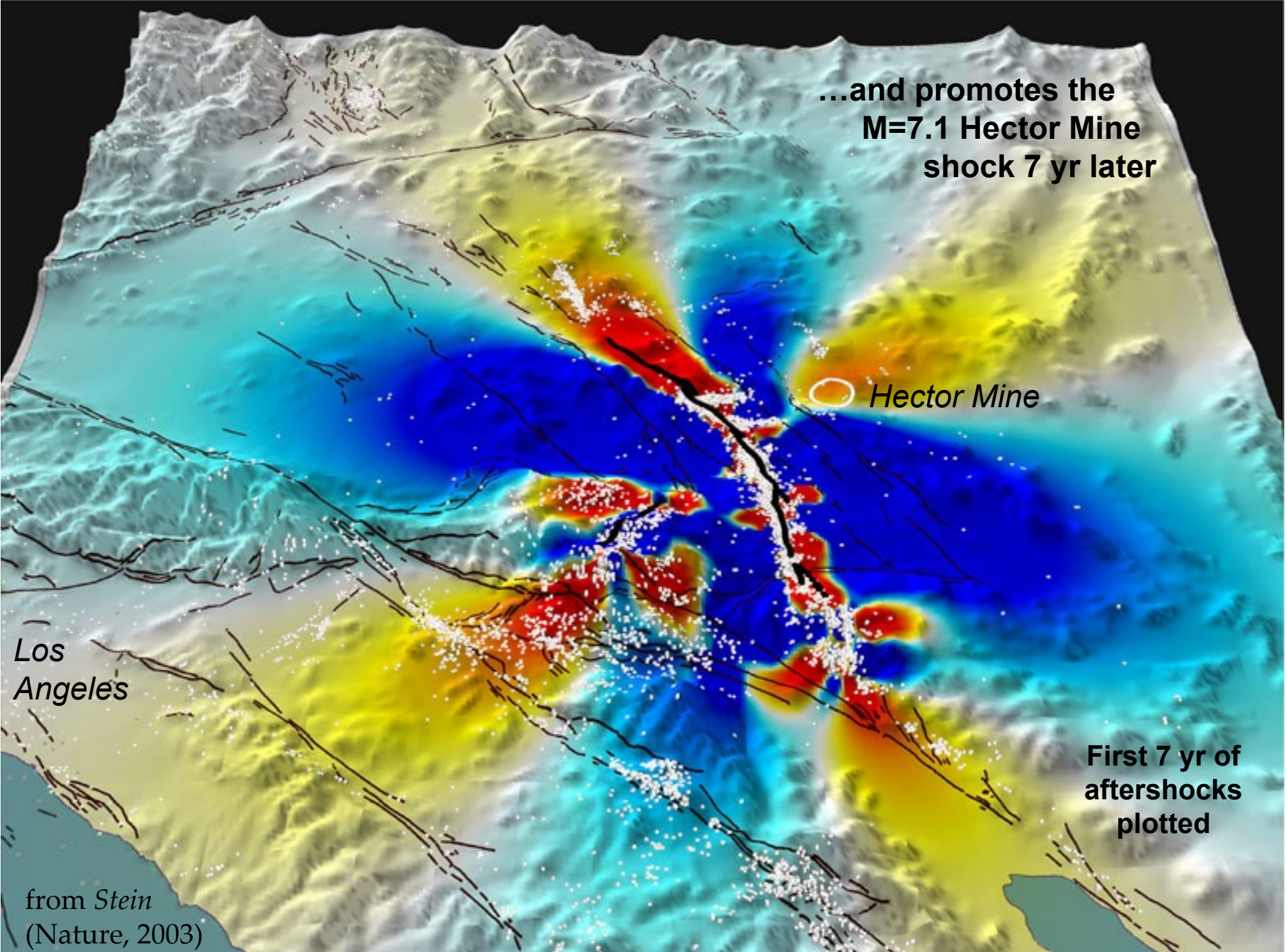
Landers

*Big
Bear*

*Los
Angeles*

**First 3 hr of
Landers
aftershocks
plotted**

*from Stein
(Nature, 2003)*



...and promotes the
M=7.1 Hector Mine
shock 7 yr later

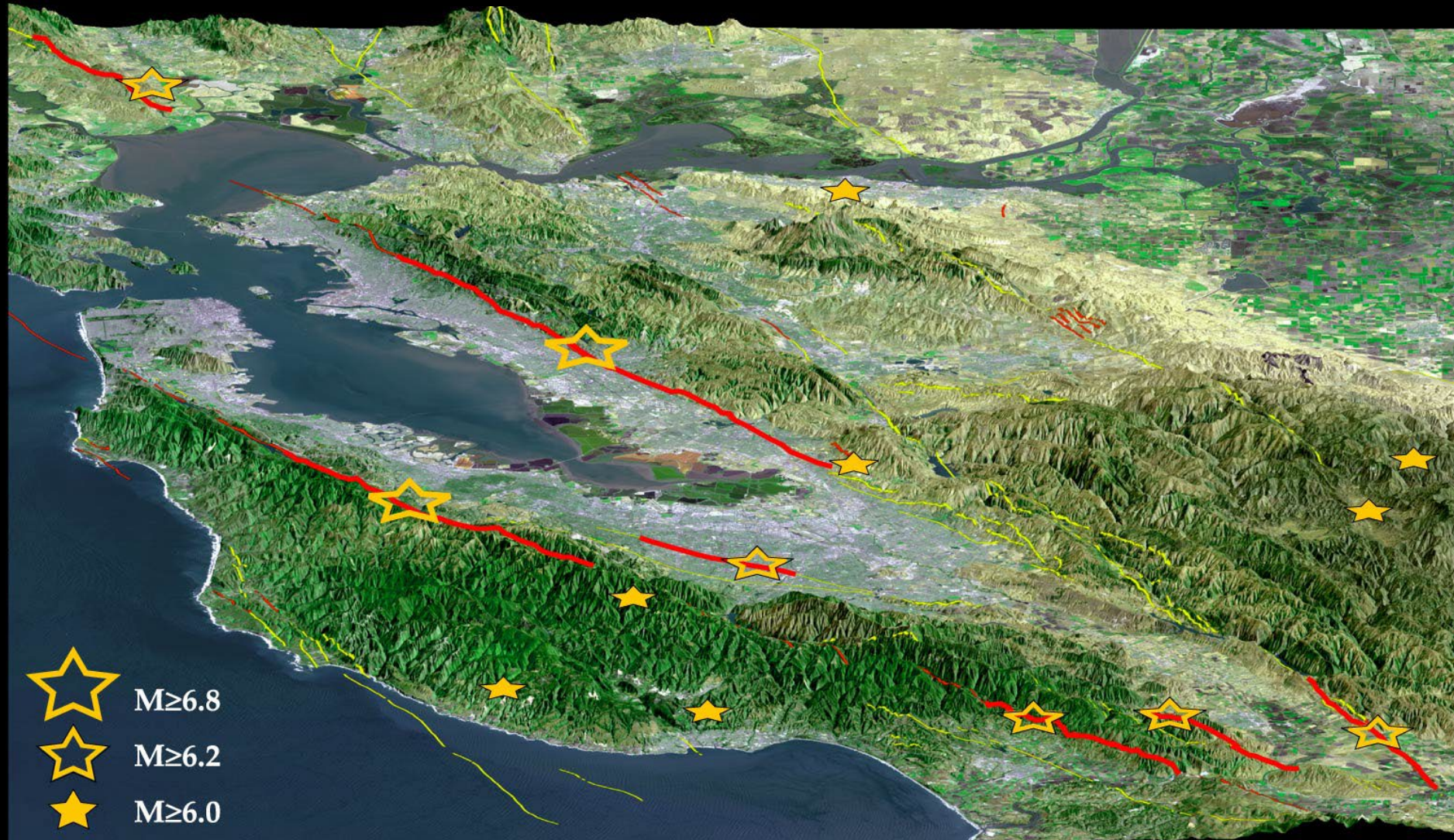
Hector Mine

Los
Angeles

First 7 yr of
aftershocks
plotted

from Stein
(Nature, 2003)

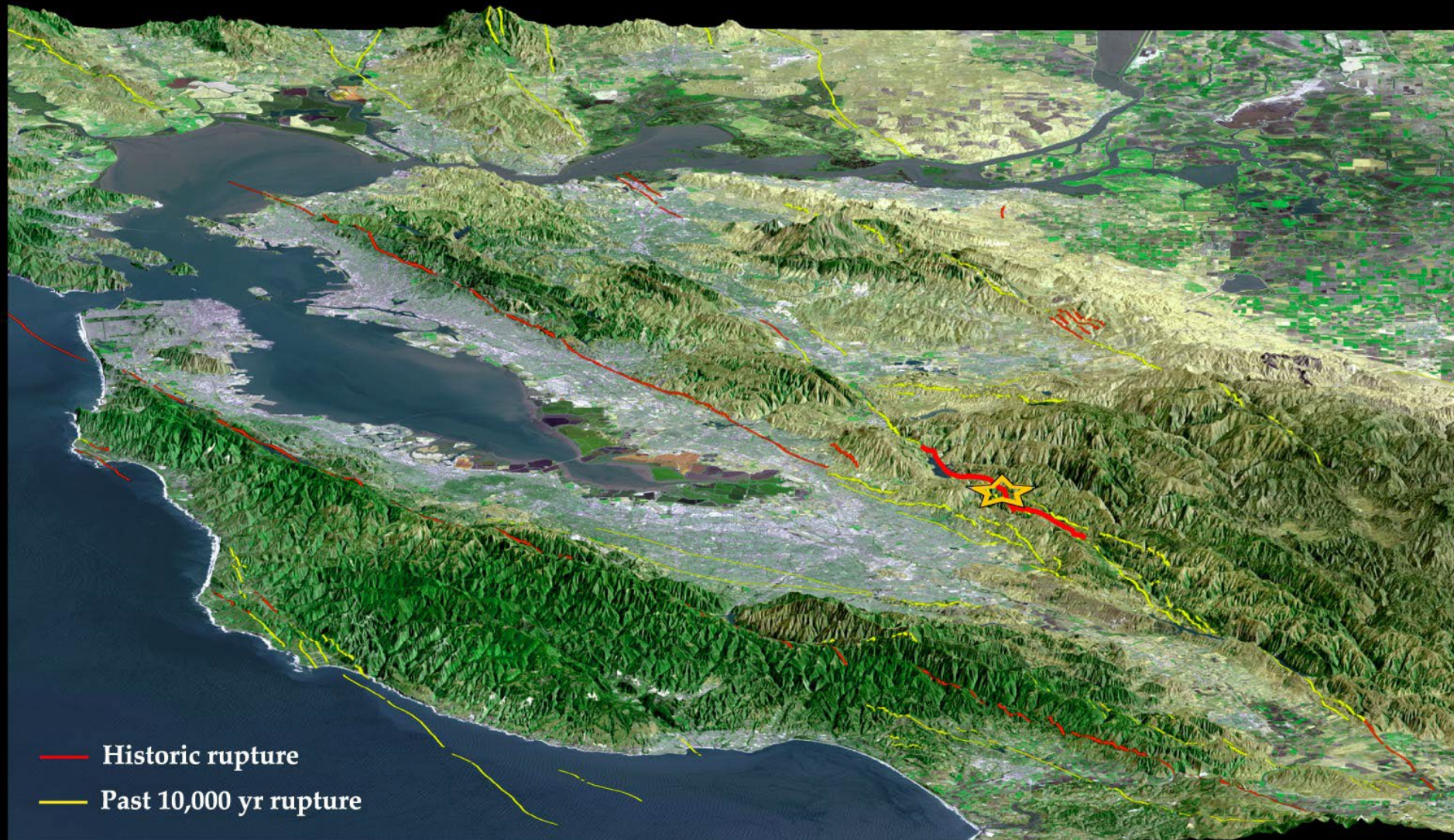
Bay area shocks during the 75 years *before* 1906



from *Stein* (Nature, 2003)

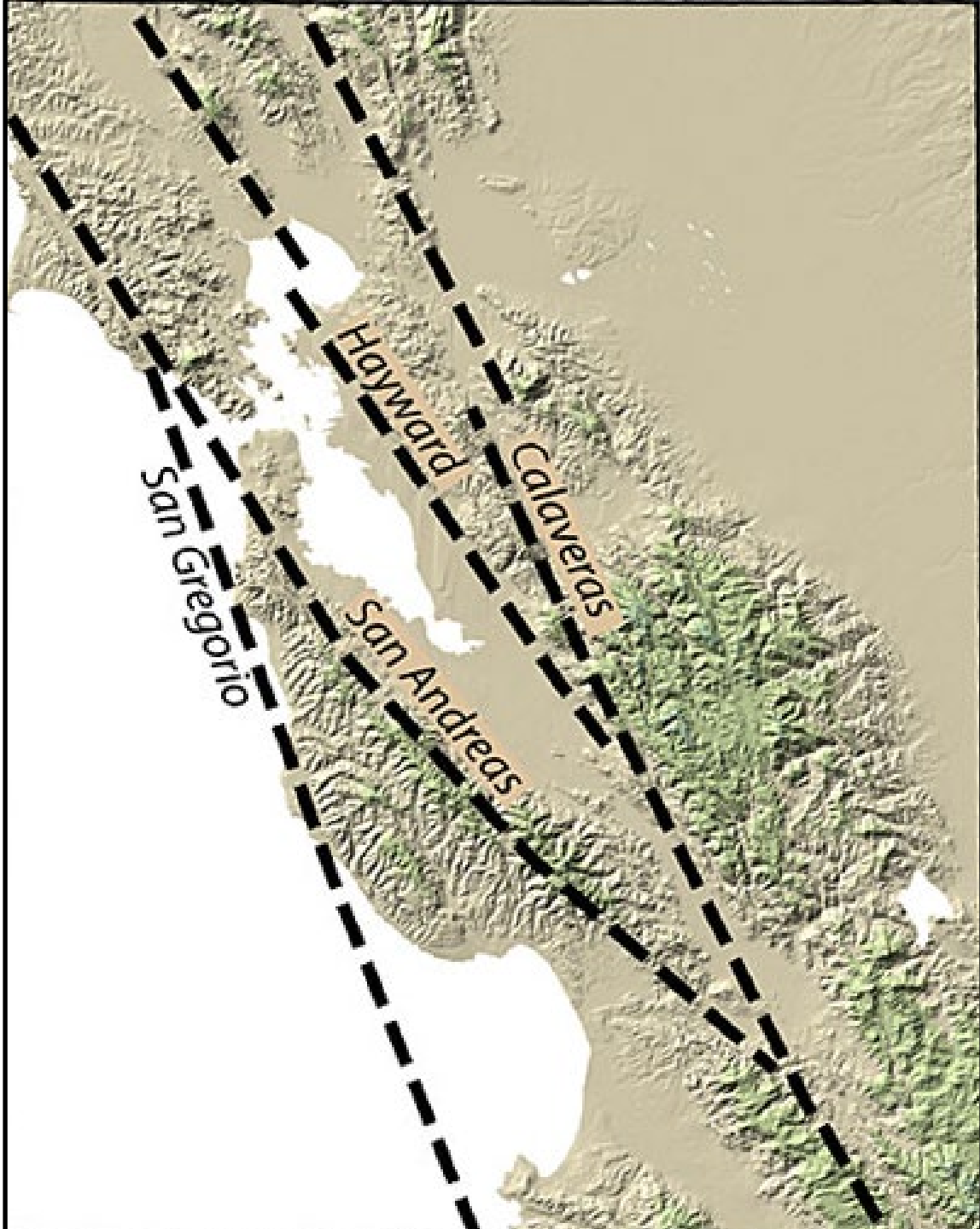
Earthquakes from *Bakun* [1999] and *Ellsworth* [1990]

Bay area shocks during the 75 years *after* 1906

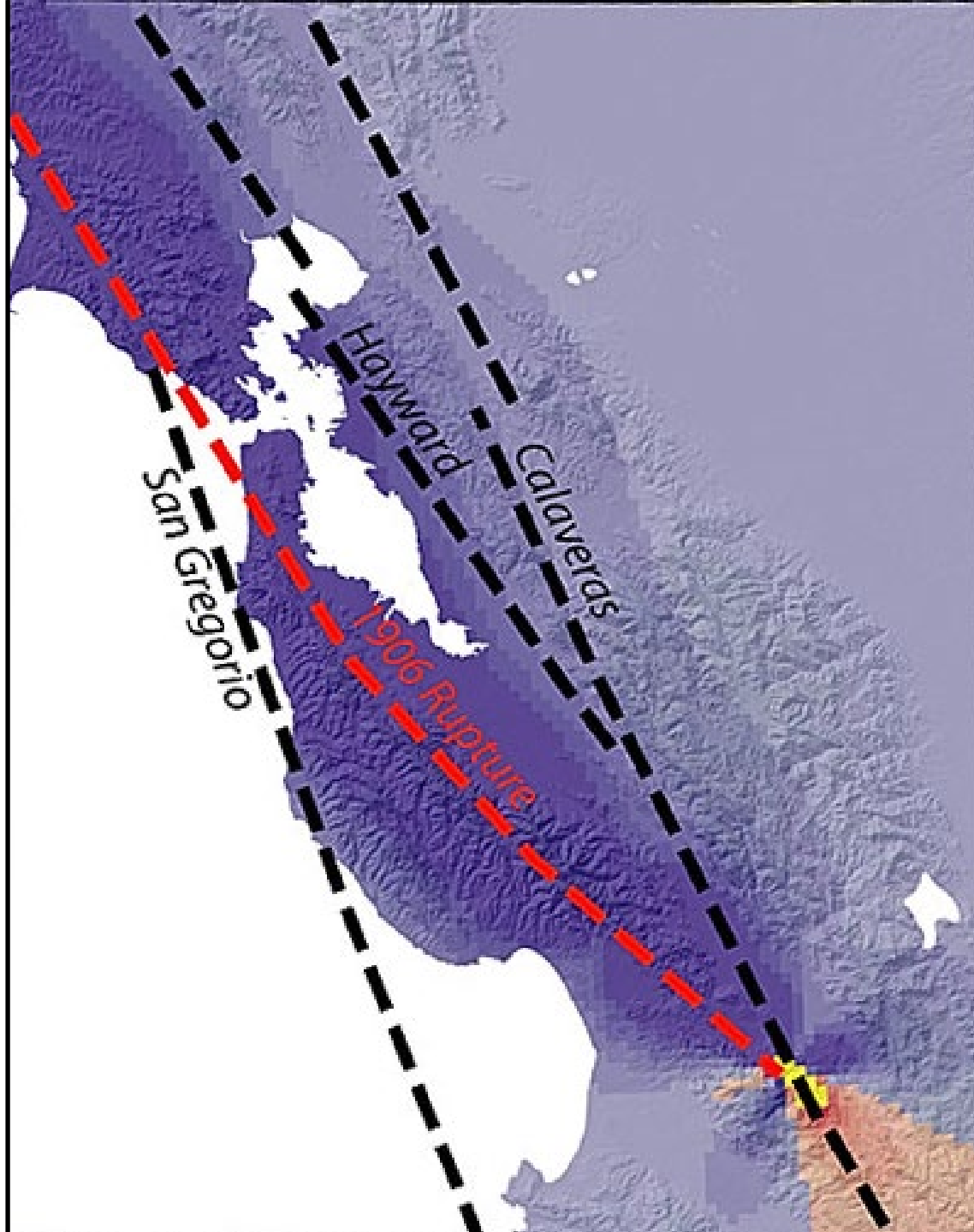


from *Stein* (Nature, 2003)

1911 M=6.2 shock from *Bakun* [BSSA, 1999]

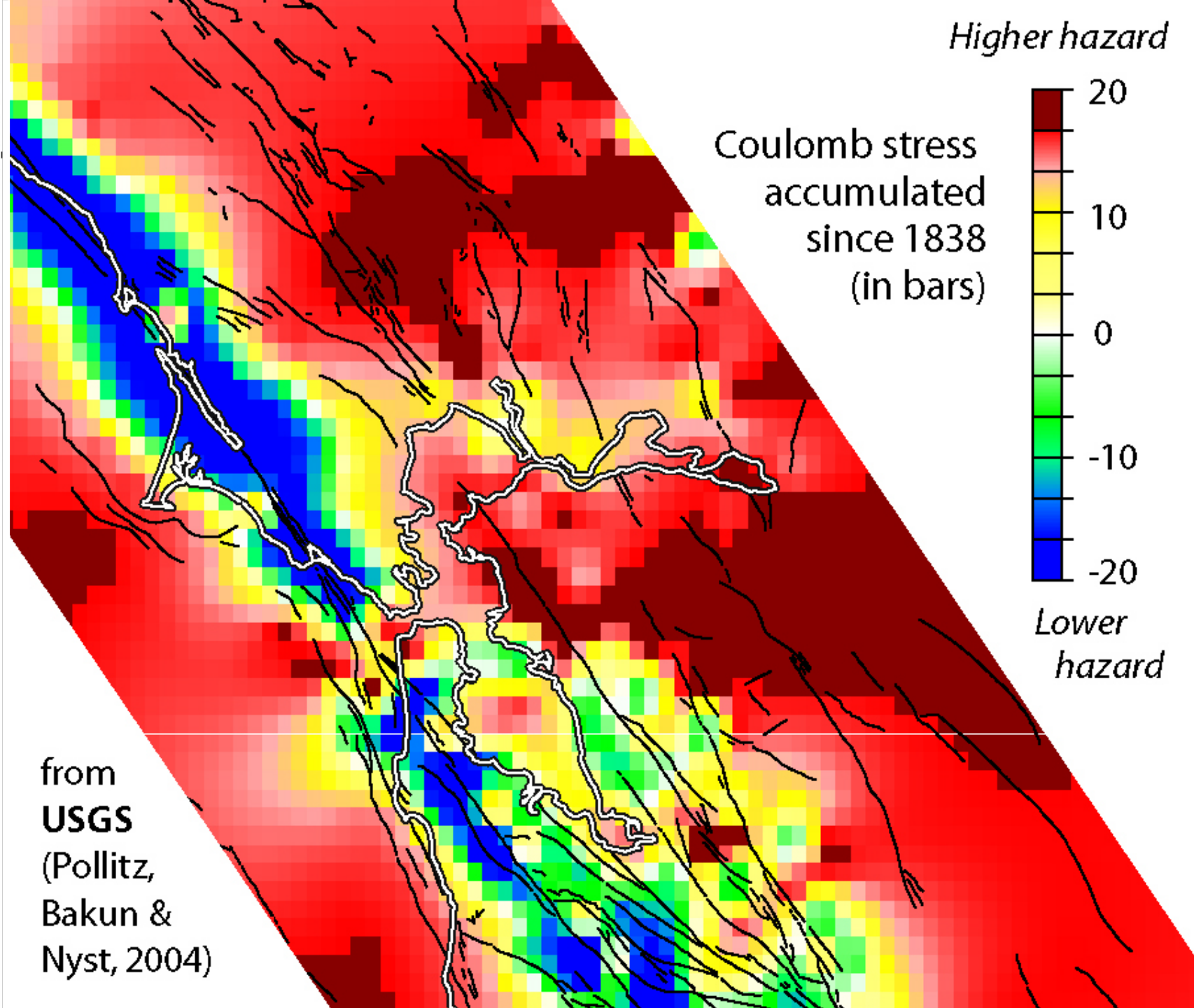


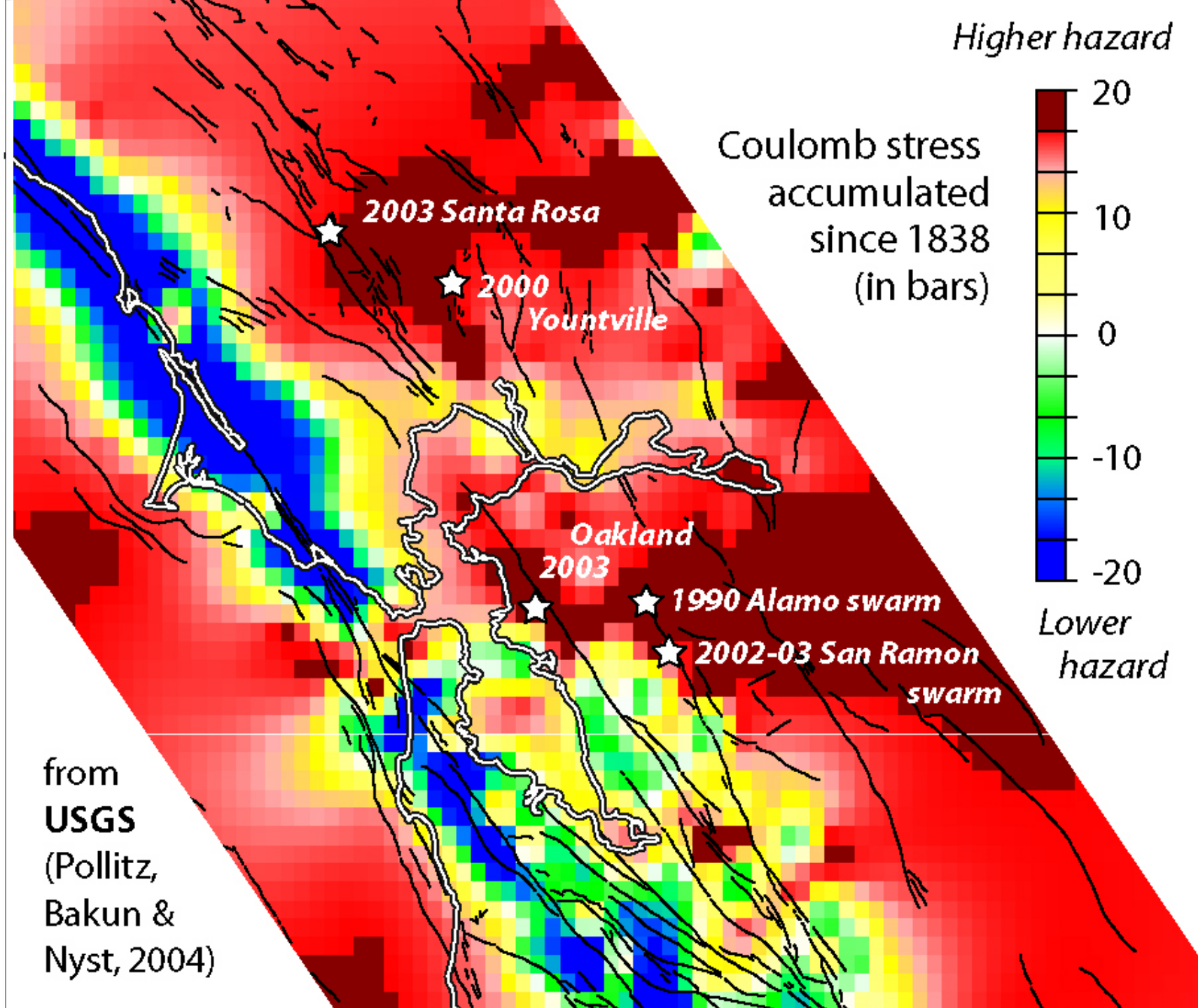
Bay area is
a system of
roughly
parallel
faults



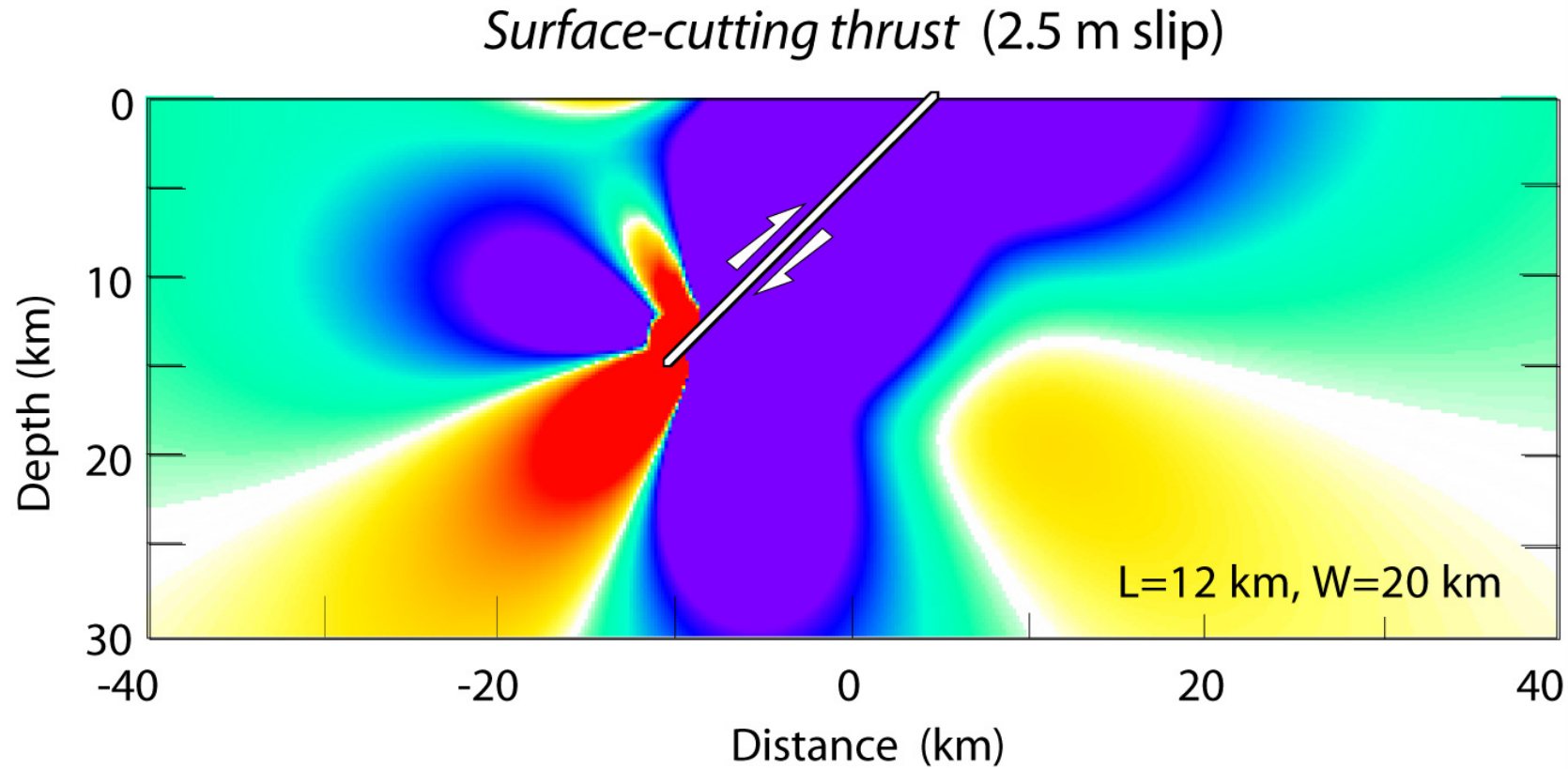
Bay area faults
may have
fallen under
a stress shadow
in 1906

from
Harris & Simpson
(1998) and *Parsons*
(2003)

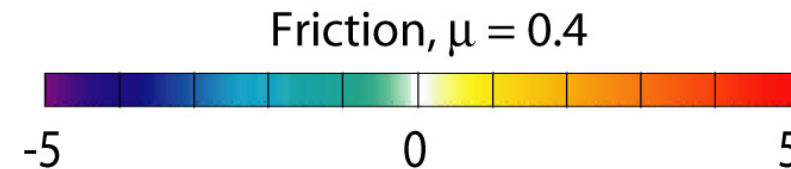




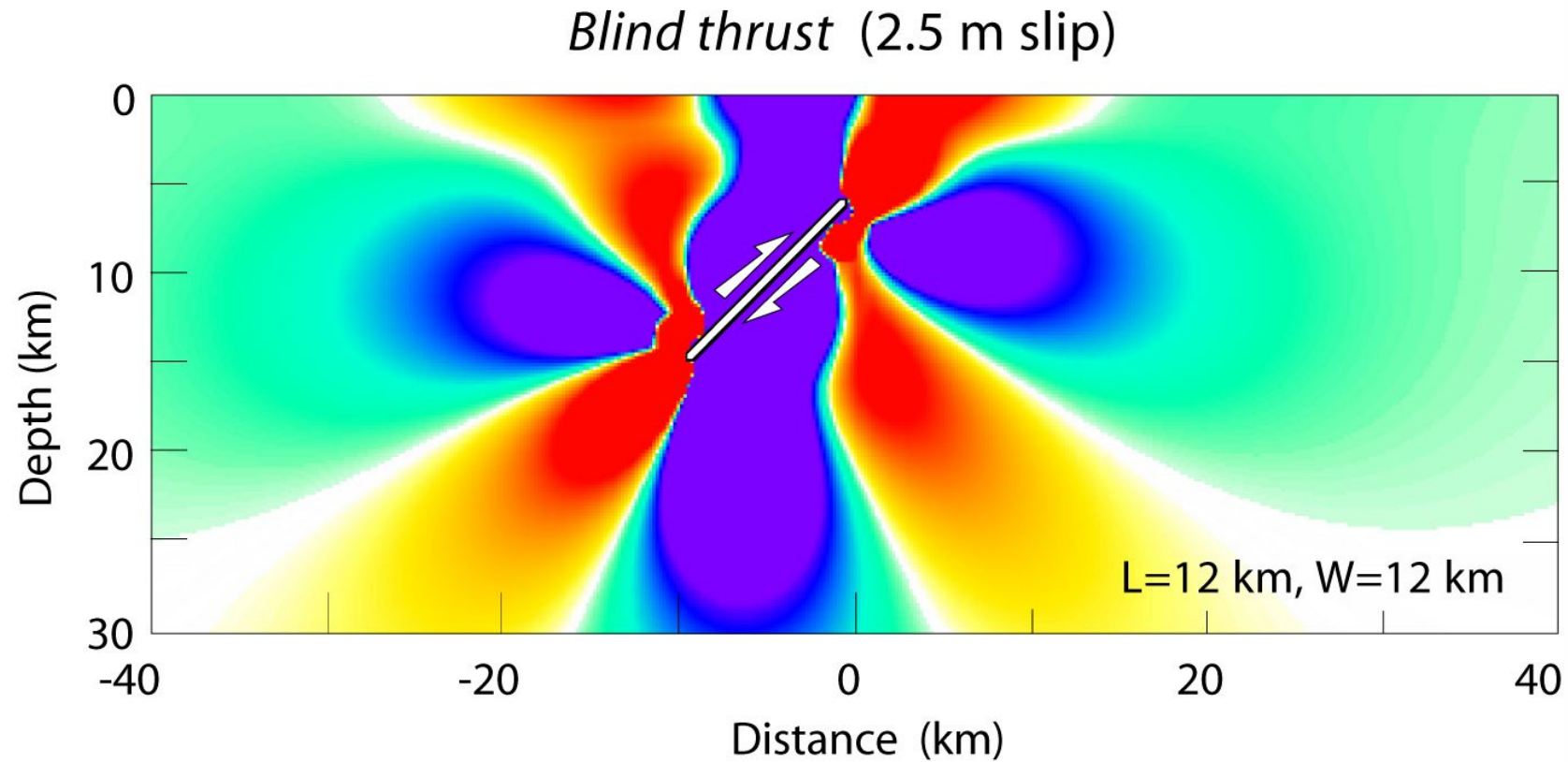
Surface-cutting thrusts drop the stress in the upper crust



Coulomb stress change (bars) on
optimally oriented thrust faults

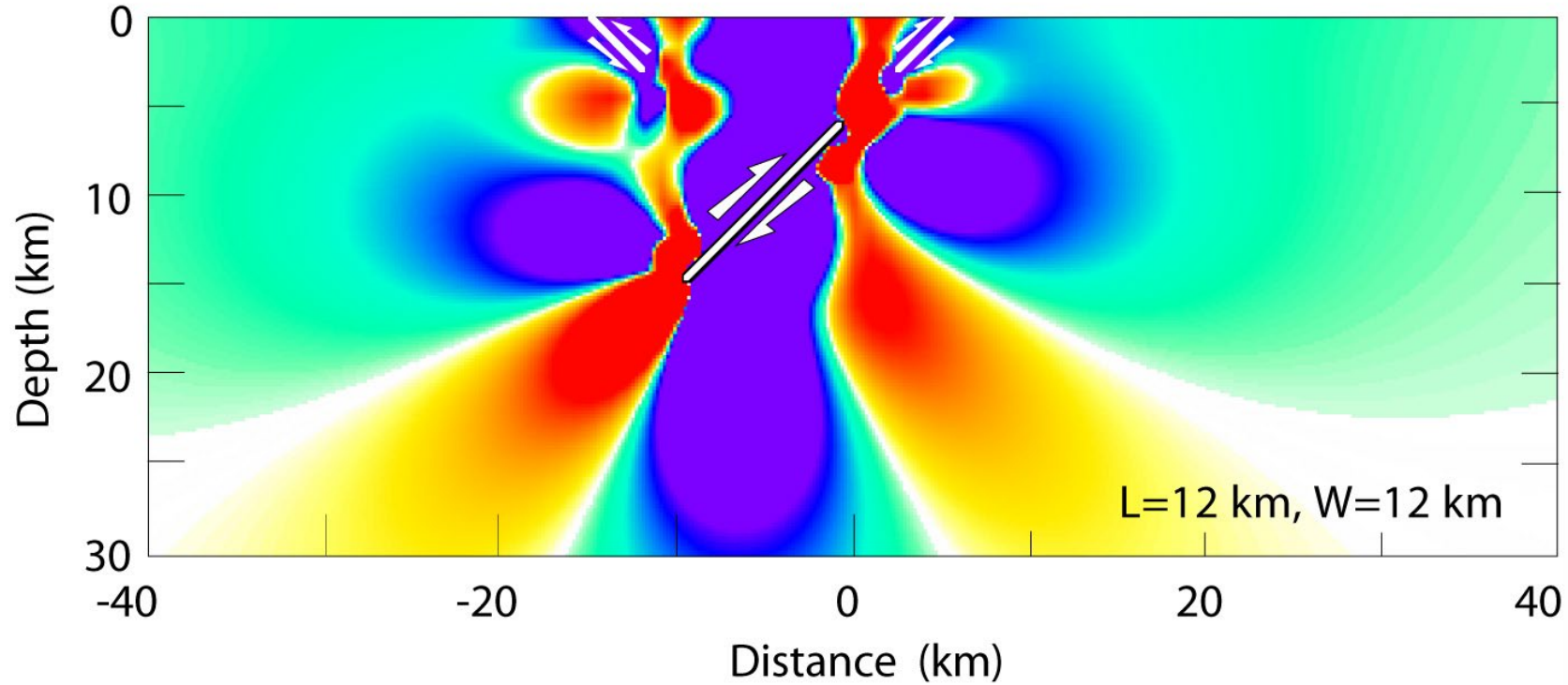


Blind thrusts raise the stress in parts of the upper crust



Secondary surface faults relieve the imparted stress

Blind thrust and secondary surface faults (each with 0.4 m slip)

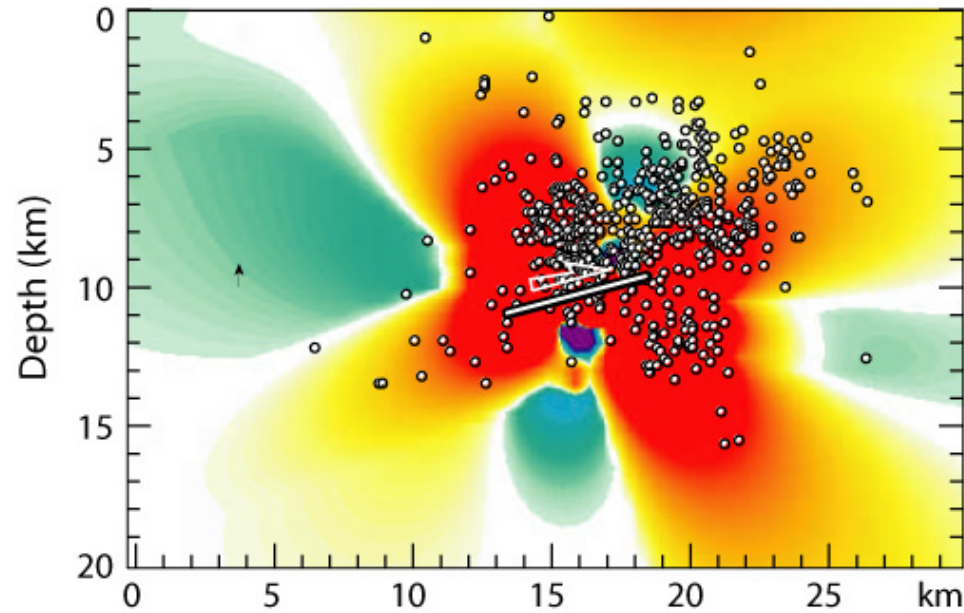


Coulomb stress change (bars) on
optimally oriented thrust faults



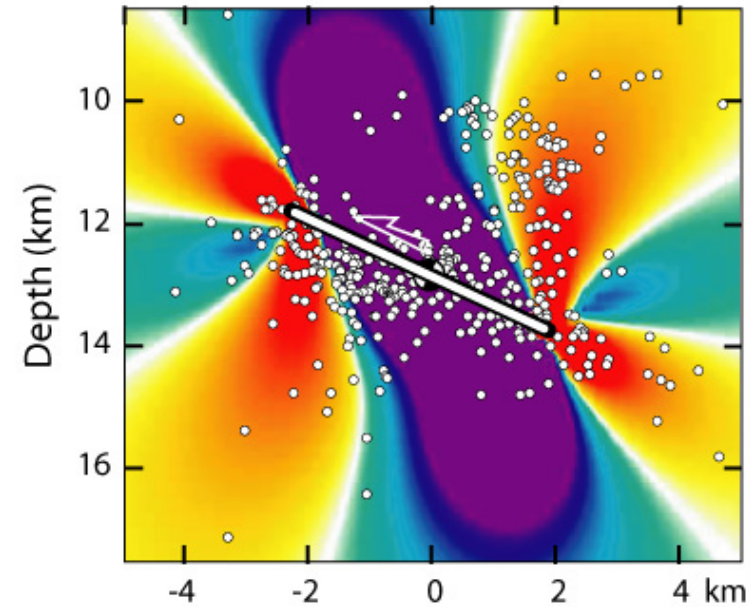
Diffuse aftershocks are characteristic of blind thrust events

1983 M=6.7 Coalinga earthquake



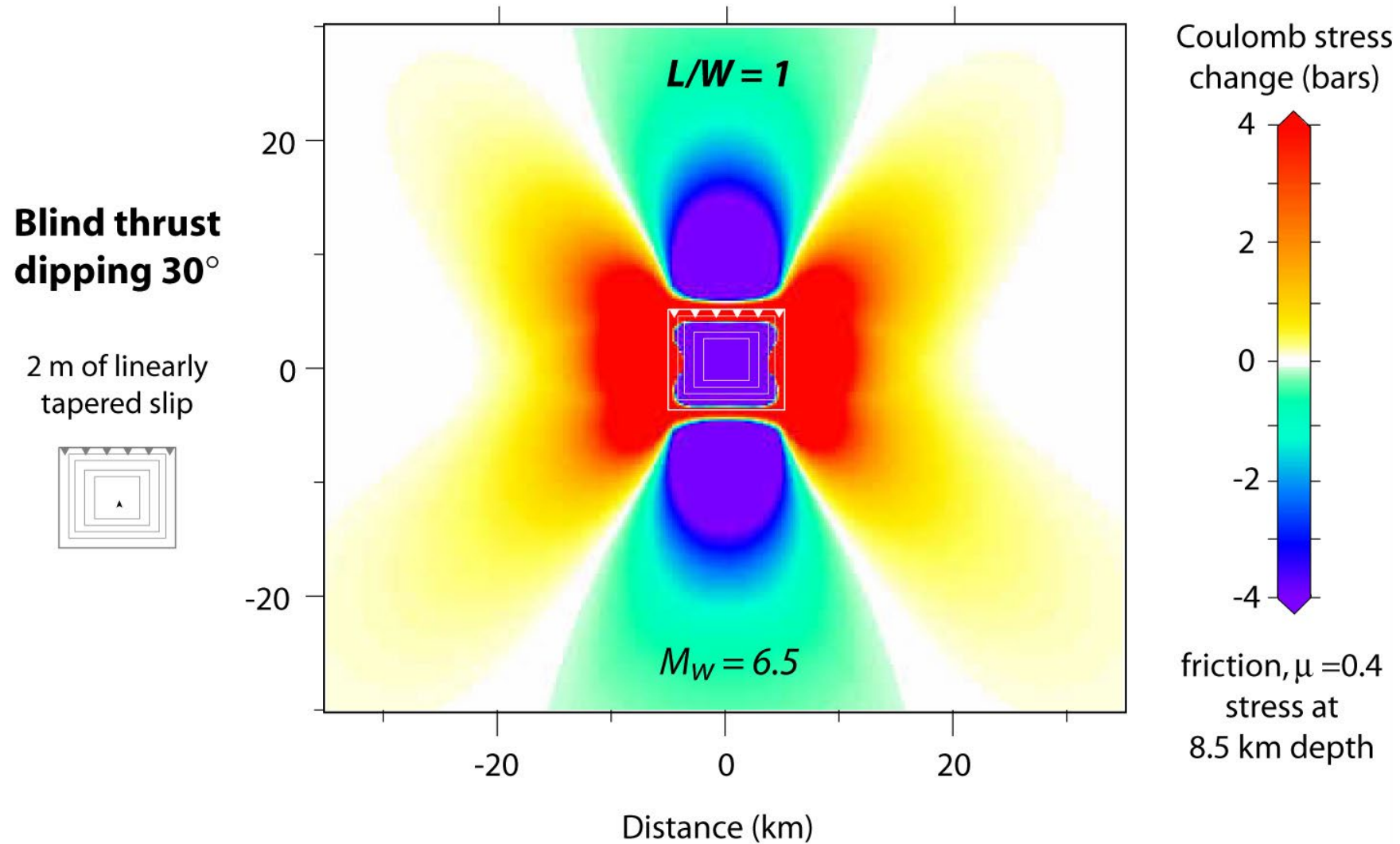
Coulomb Stress Change on optimally oriented thrust fault faults

1987 M=6.0 WhittierNarrows quake



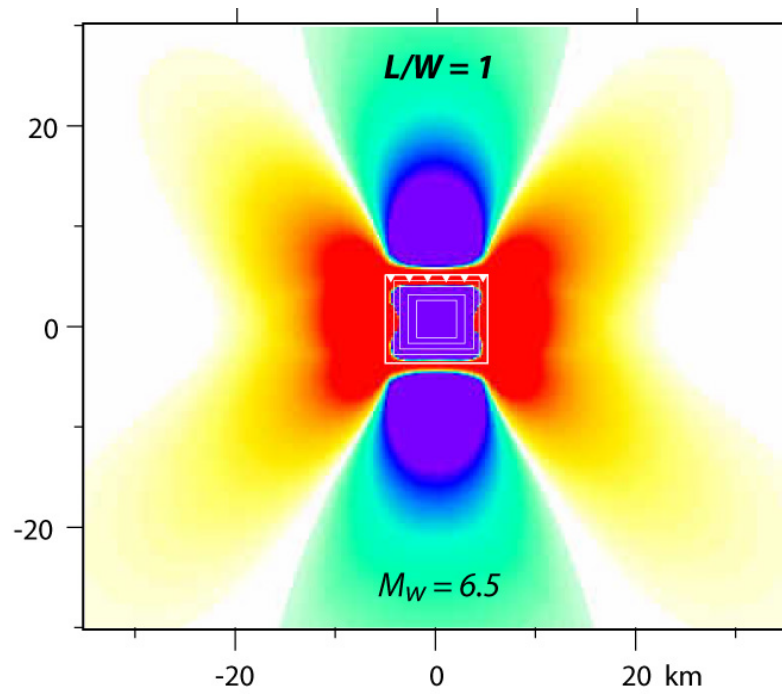
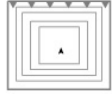
Coulomb stress change on thrust faults parallel to the source fault rupture

Earthquake increases Coulomb stress in 'butterfly wings'

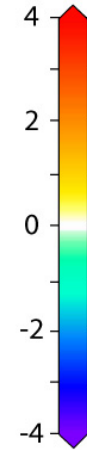


Blind thrust dipping 30°

2 m of linearly tapered slip



Coulomb stress change (bars)



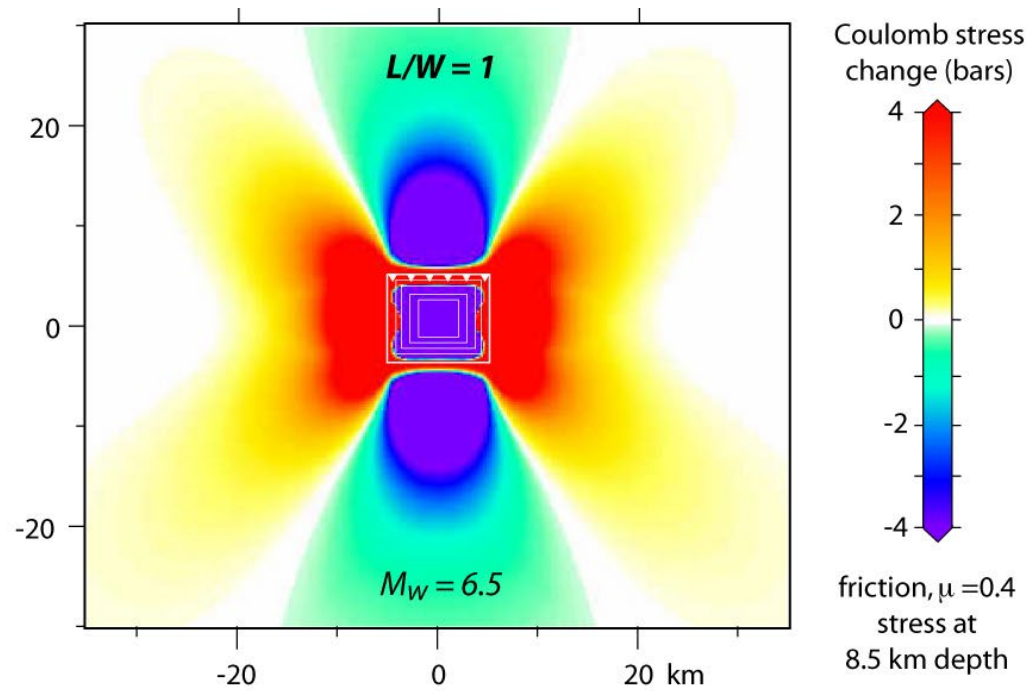
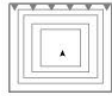
friction, $\mu = 0.4$
stress at
8.5 km depth

Stress is imparted to the 'butterfly wings'

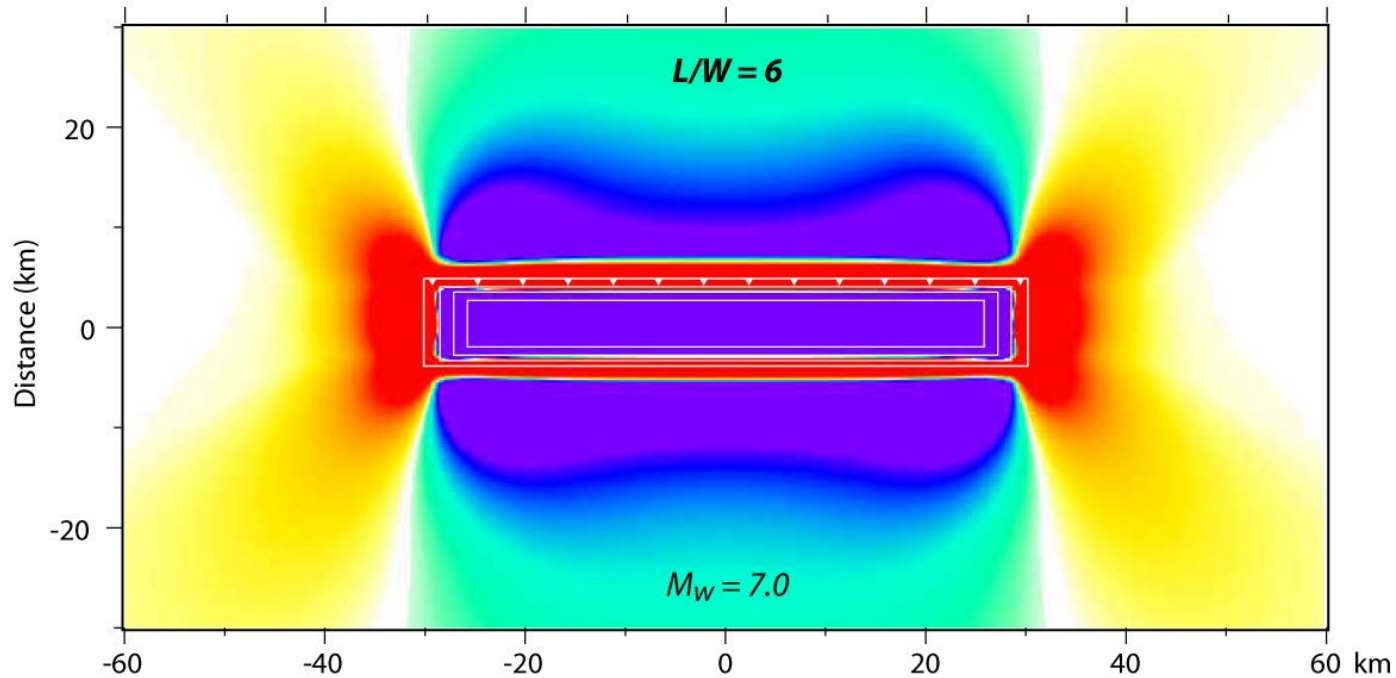
from *Lin & Stein*
(JGR, 2004)

Blind thrust dipping 30°

2 m of linearly tapered slip



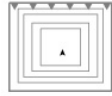
Short thrust
much more
efficient at
transferring
stress
along strike



from *Lin & Stein*
(JGR, 2004)

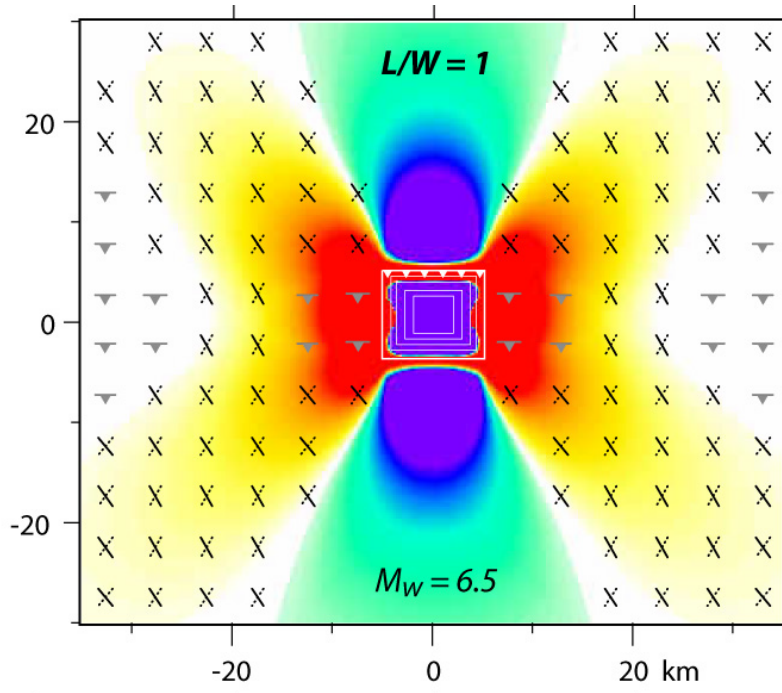
Blind thrust dipping 30°

2 m of linearly tapered slip

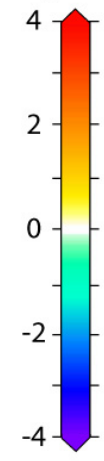


Optimum receiver planes

- ▼ thrust
- ⊗ left-lateral
- ⊗ right-lateral

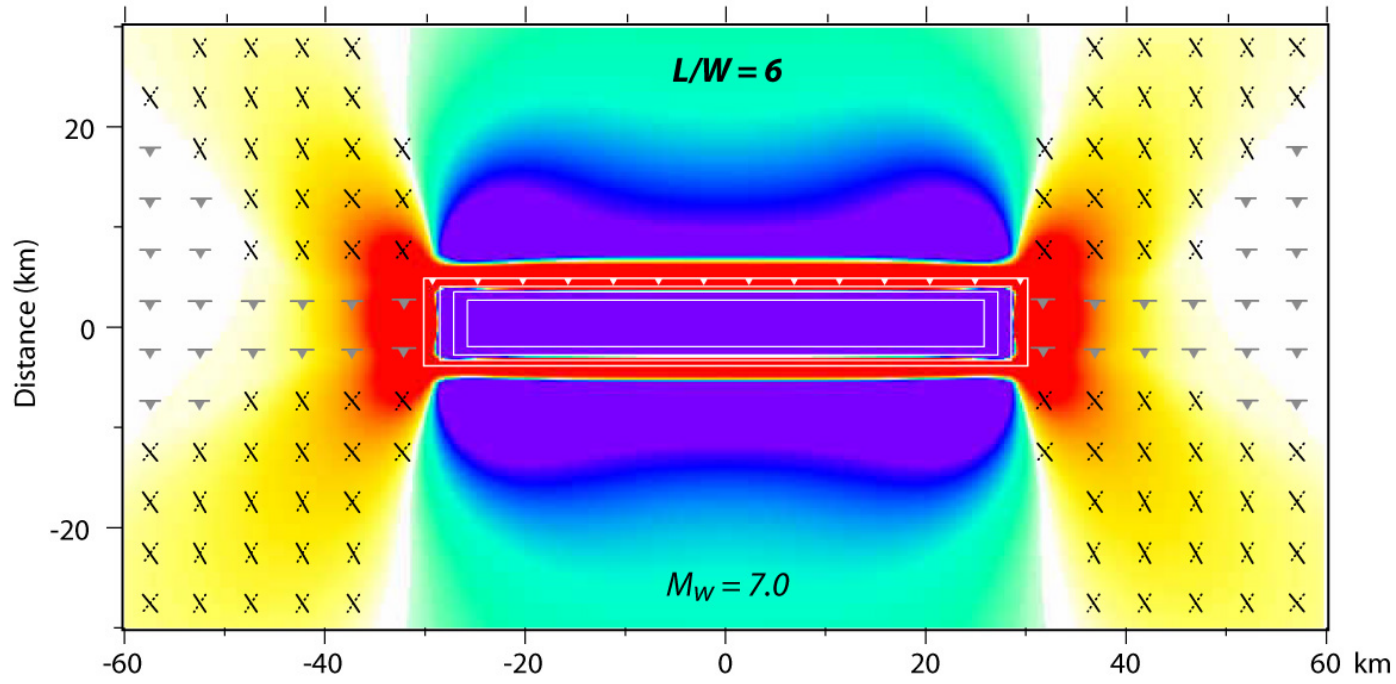


Coulomb stress change (bars)



friction, $\mu = 0.4$
stress at
8.5 km depth

Stress imparted to wings promotes strike-slip faulting



from *Lin & Stein*
(JGR, 2004)

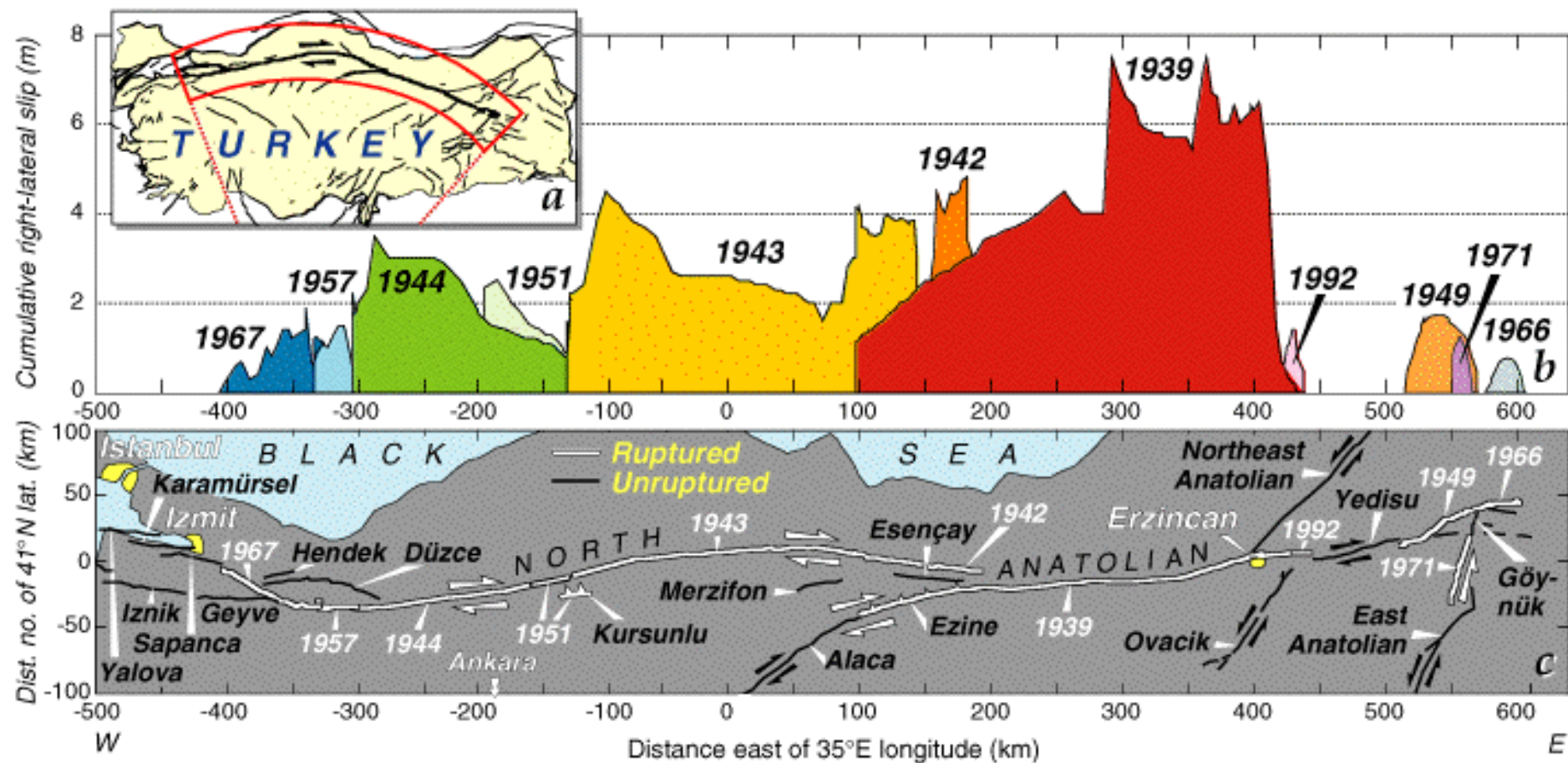


Figure 1 17 Oct 96 Stein et al.

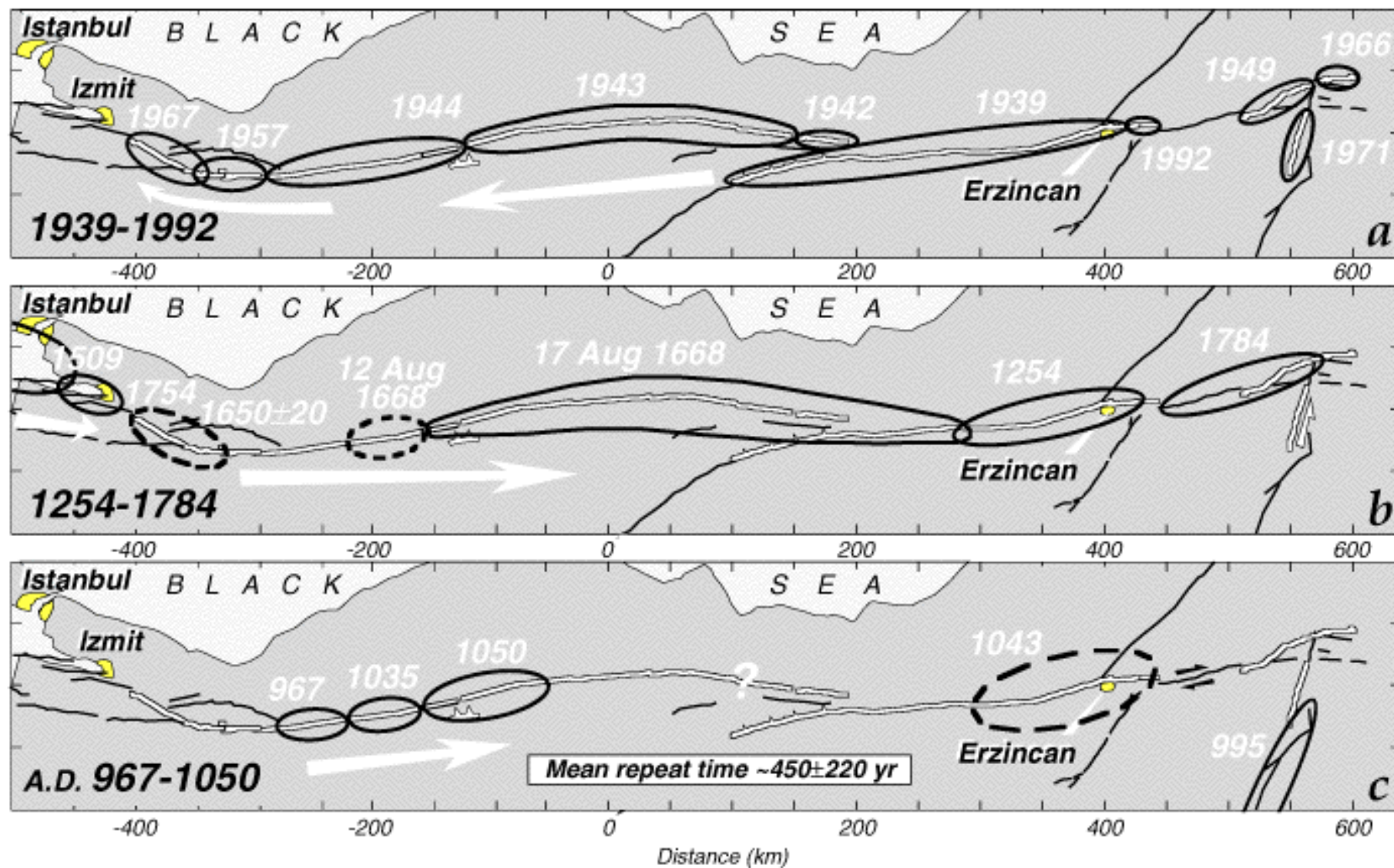
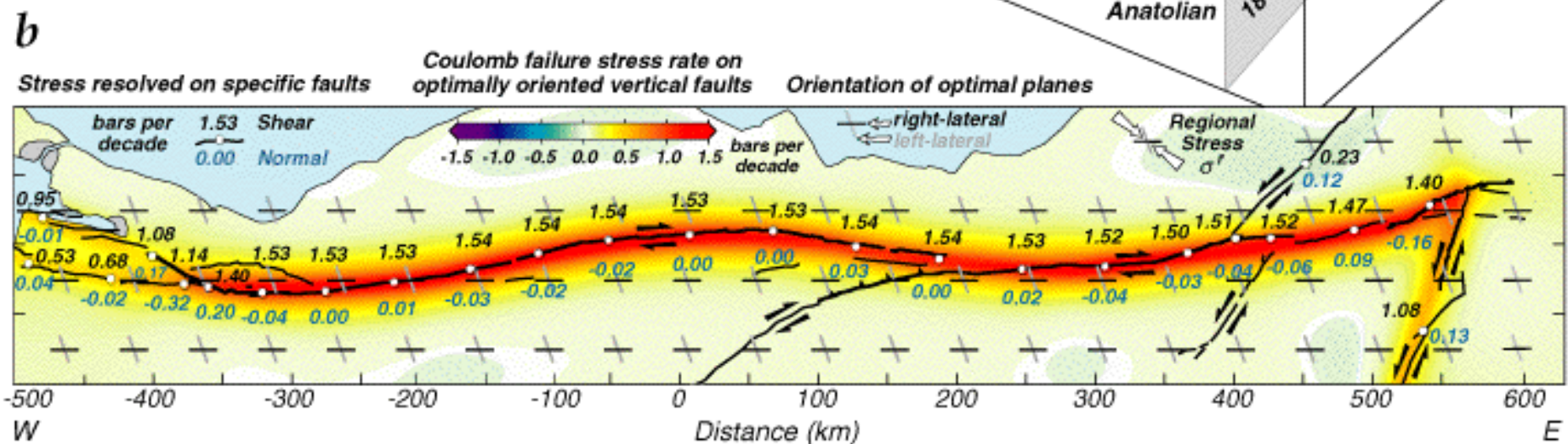
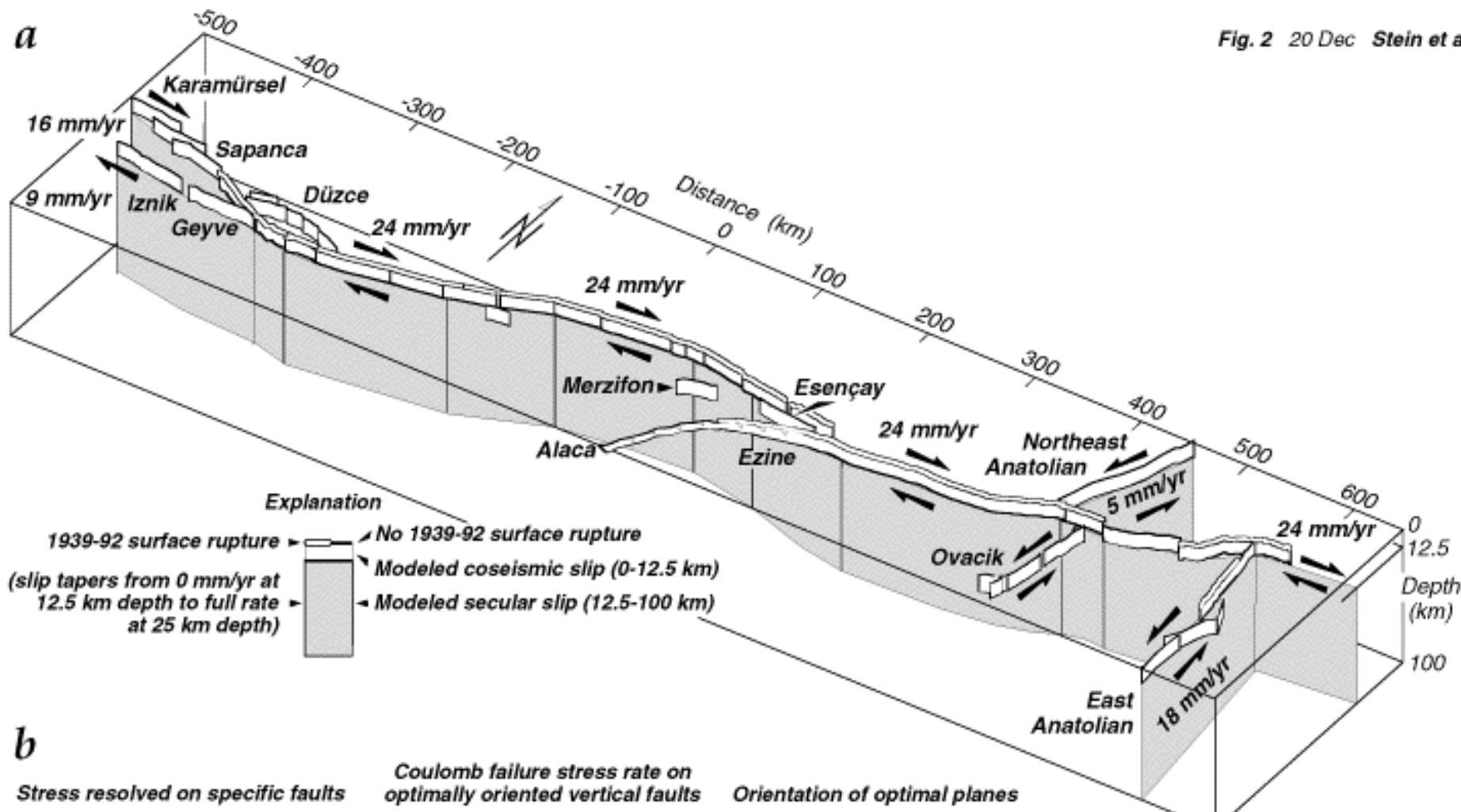


Fig. 7 17 Oct 96 Stein et al.



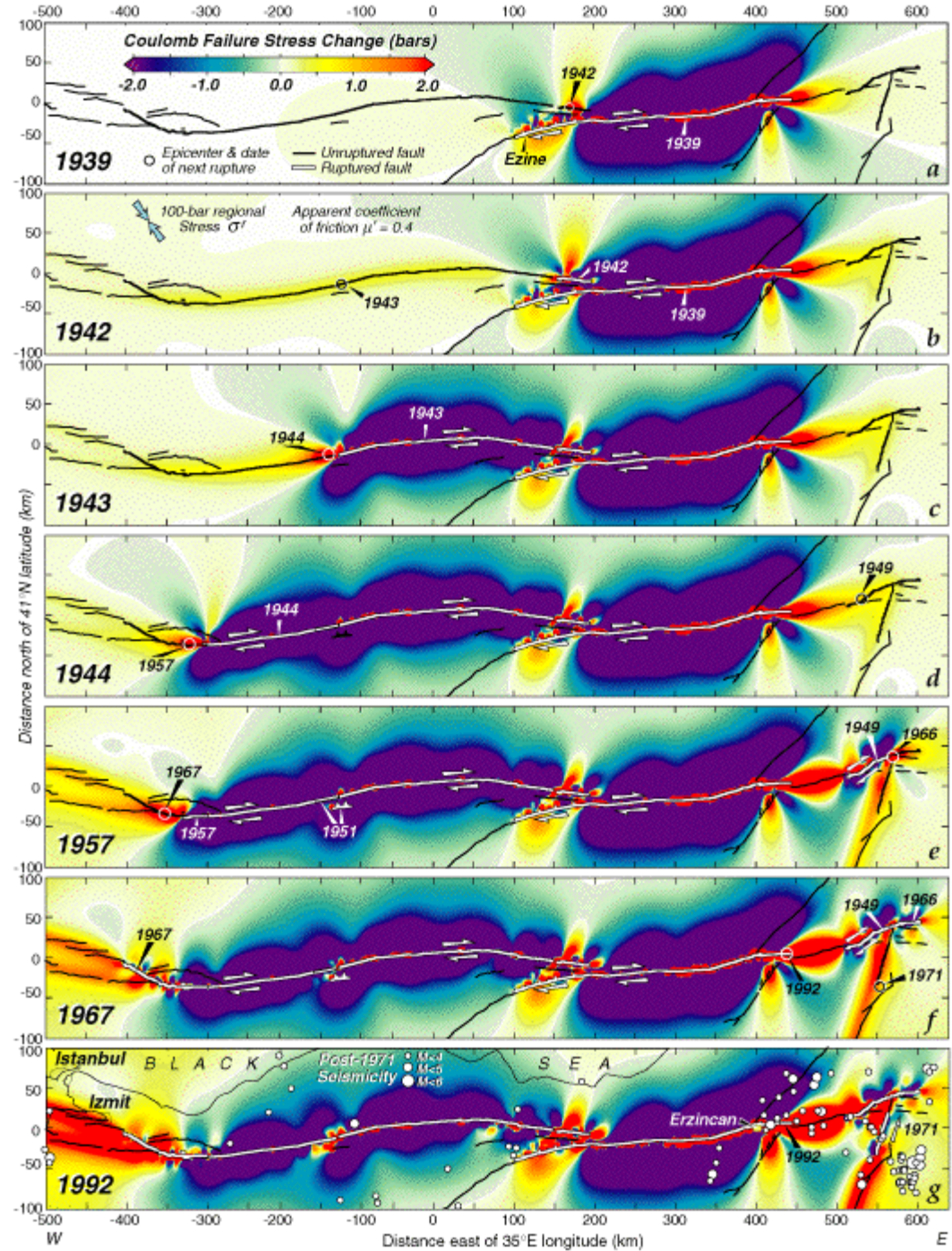


Figure 4 17 Oct 96 Stein et al.

From Stress Change to Earthquake Probability Change

